

Squeezing with parametric modulation close to quantum regime

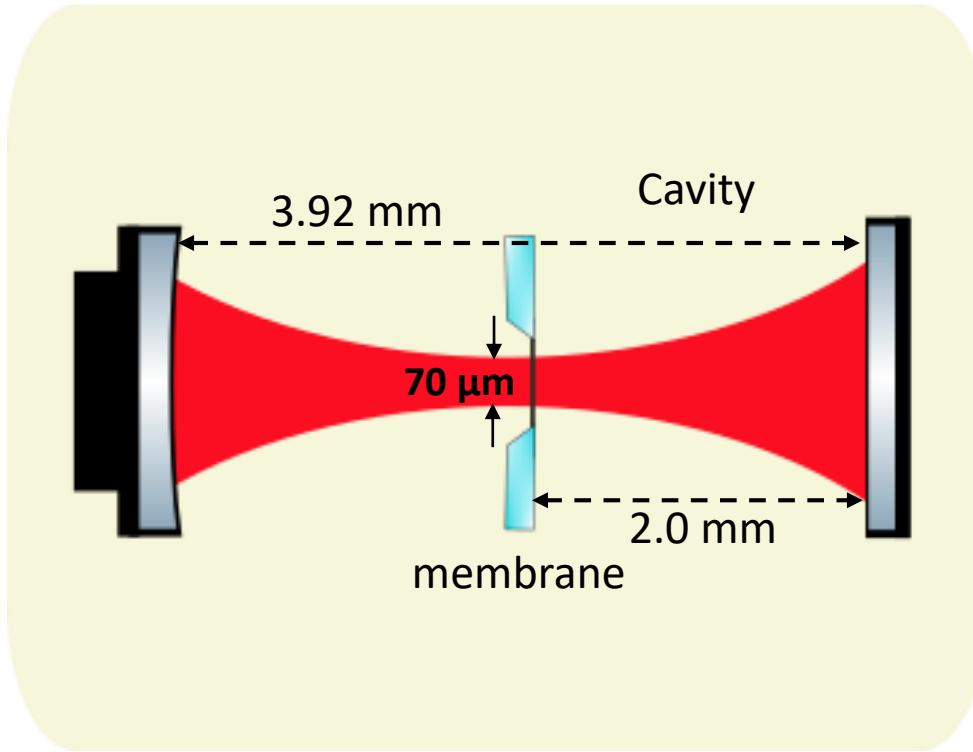
Avishek Chowdhury

Paolo Vezio

Francesco Marino

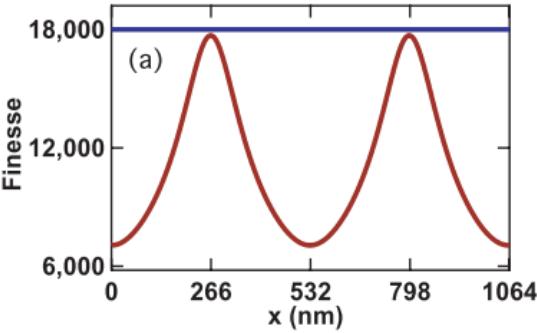
Francesco Marin

Cavity opto-mechanics: membrane in the middle

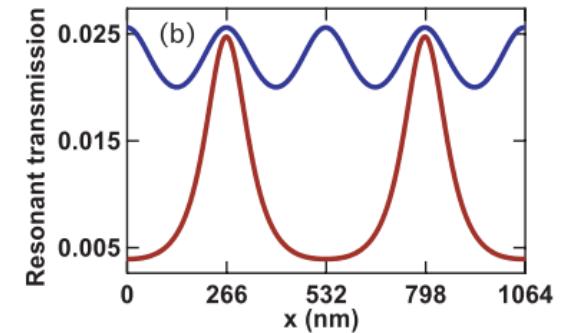


Cavity:

- Finesse $\sim 20,000$
- Cavity linewidth $\sim 1.9 \text{ MHz}$



Jayich et al. NJP 10 (2008) 095008



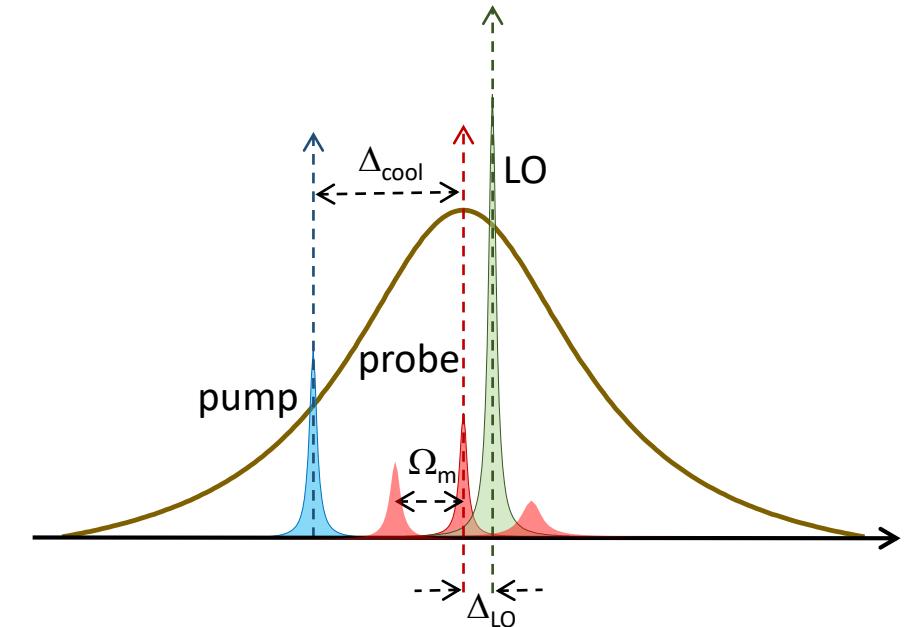
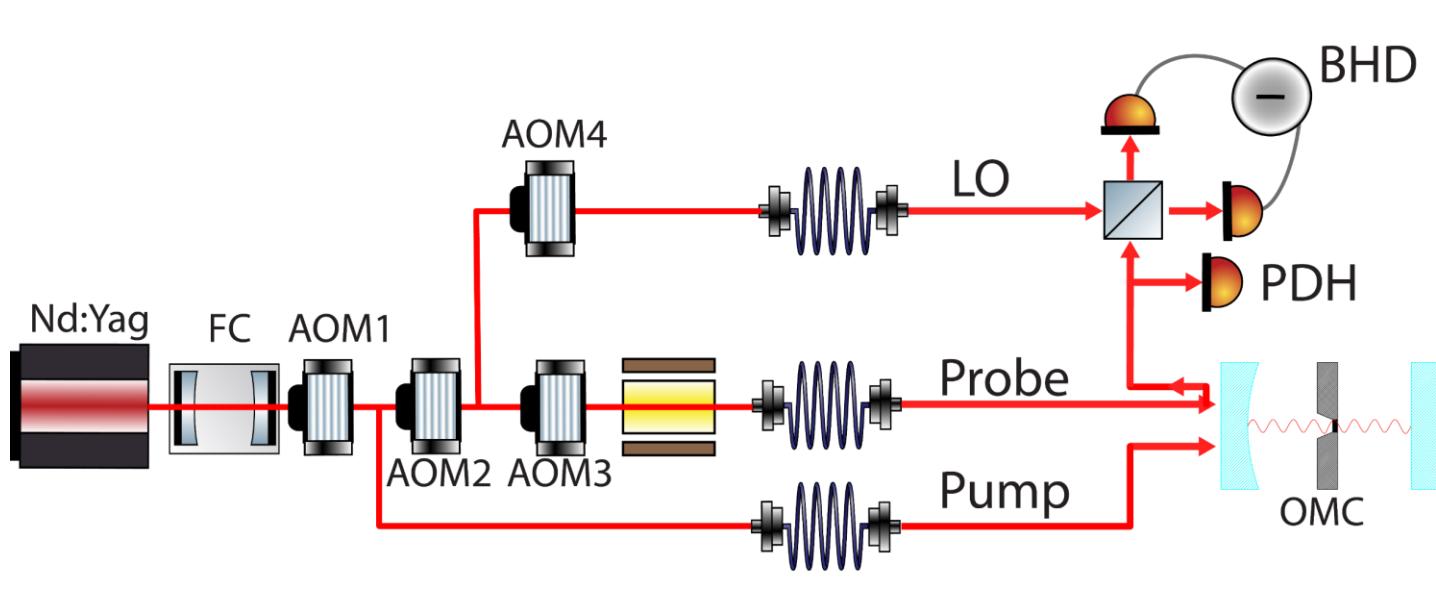
Tunable finesse as a function of the membrane position



Membrane:

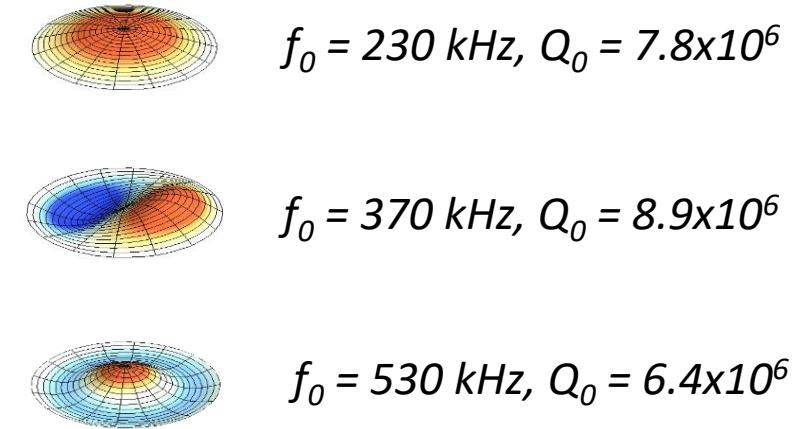
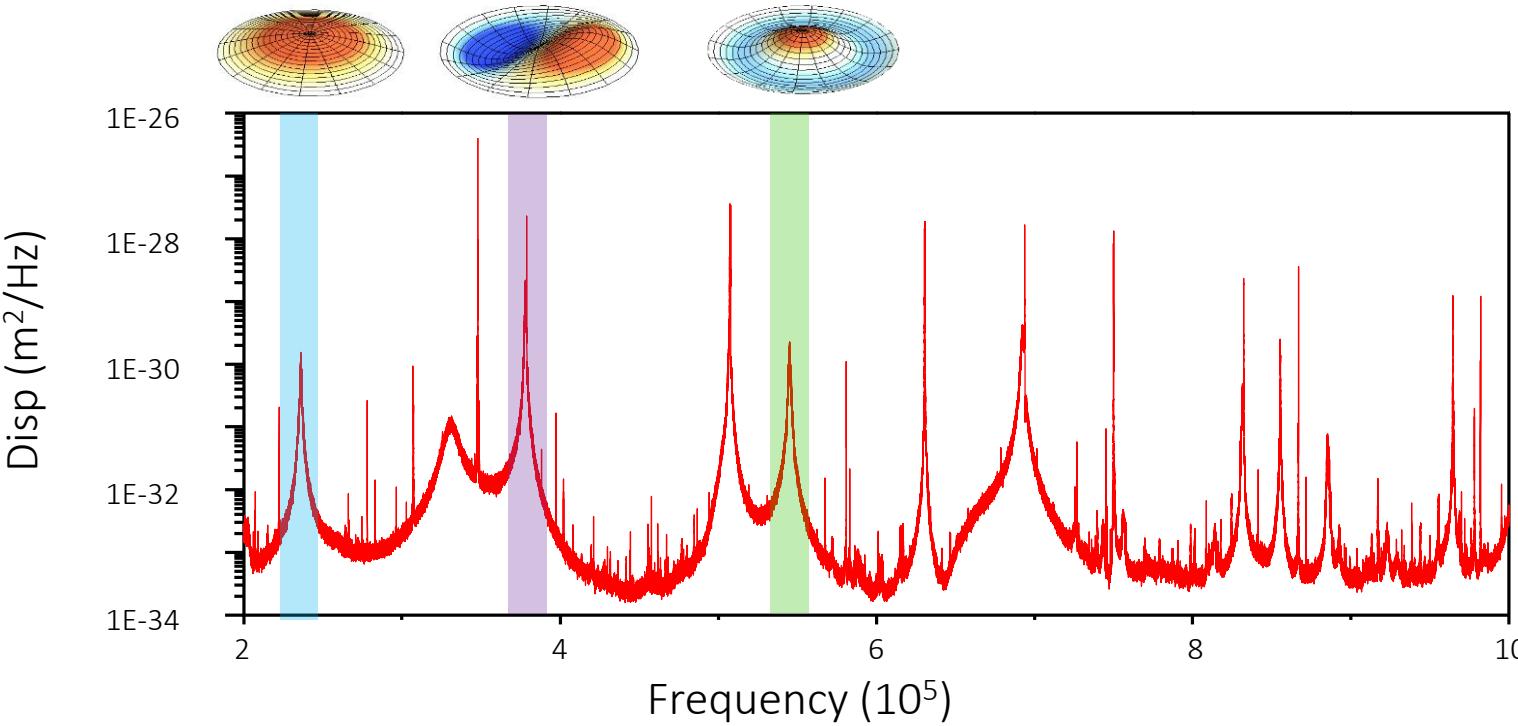
- Thickness $\sim 100 \text{ nm}$
- Diameter $\sim 1.64 \text{ mm}$
- $Q_{\text{cryogenic temp}} \sim 10^7$

Simplified optical set-up



- **Weak probe**: to detect the displacement of the mechanical mode, in resonance with the cavity.
- **Strong pump**: detuned from the cavity, used to optically cool down the mechanical mode.
- **LO beam**: Local Oscillator beam used to make heterodyne measurement on the mechanical mode.

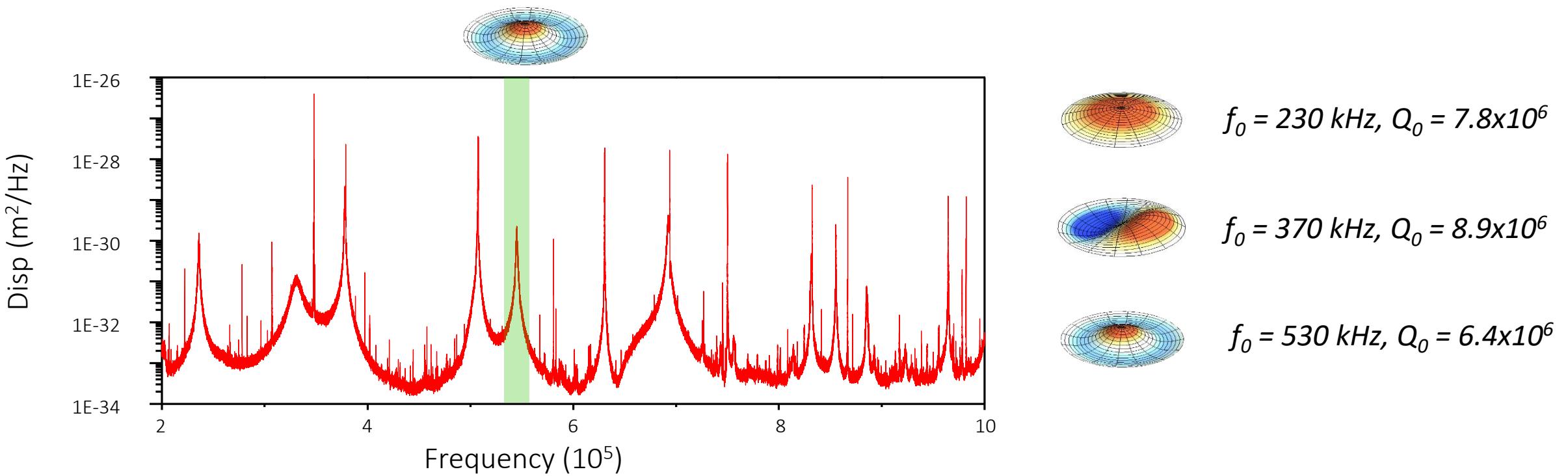
Membrane modes: eigen-modes



Choice of the modes

- Considerably high Q factor.
- High optomechanical coupling: highly dependent on the effective overlap between the laser beam and the membrane eigen mode.

Membrane modes: eigen-modes



We choose: Mode at 530 kHz

- Q factor: 6.4×10^6 .
- With an optomechanical coupling $g_0 \sim 5 \text{ Hz}$.

Cooling a mechanical mode

$$\delta\Omega_m(\omega) = g^2 \frac{\Omega_m}{\omega} \left[\frac{\Delta + \omega}{(\Delta + \omega)^2 + \kappa^2/4} + \frac{\Delta - \omega}{(\Delta - \omega)^2 + \kappa^2/4} \right]$$

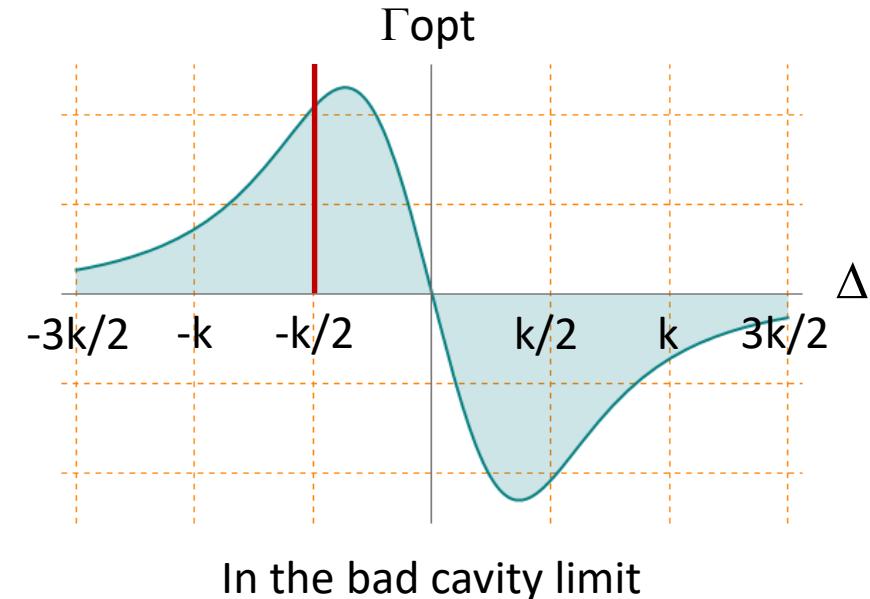
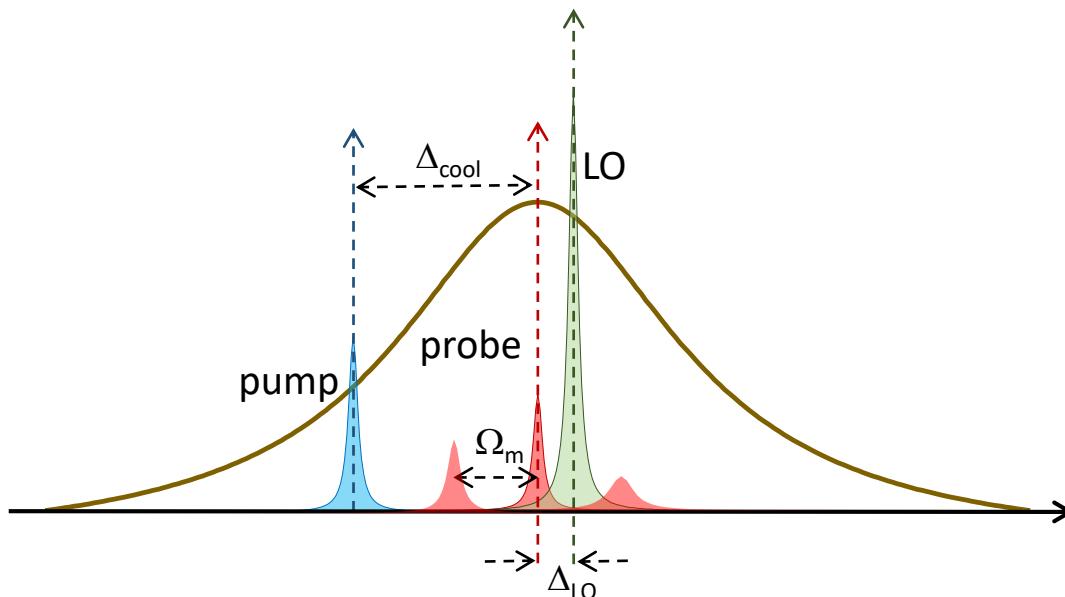
Optical Spring effect

$$\Gamma_{\text{opt}}(\omega) = g^2 \frac{\Omega_m}{\omega} \left[\frac{\kappa}{(\Delta + \omega)^2 + \kappa^2/4} - \frac{\kappa}{(\Delta - \omega)^2 + \kappa^2/4} \right]$$

Damping of the mode 'm'

$$g^2 = g_0^2 \bar{n}_{\text{cav}}$$

$$\Gamma_{\text{opt}} \sim f(n_{\text{cav}}, \Delta)$$



Cooling a mechanical mode

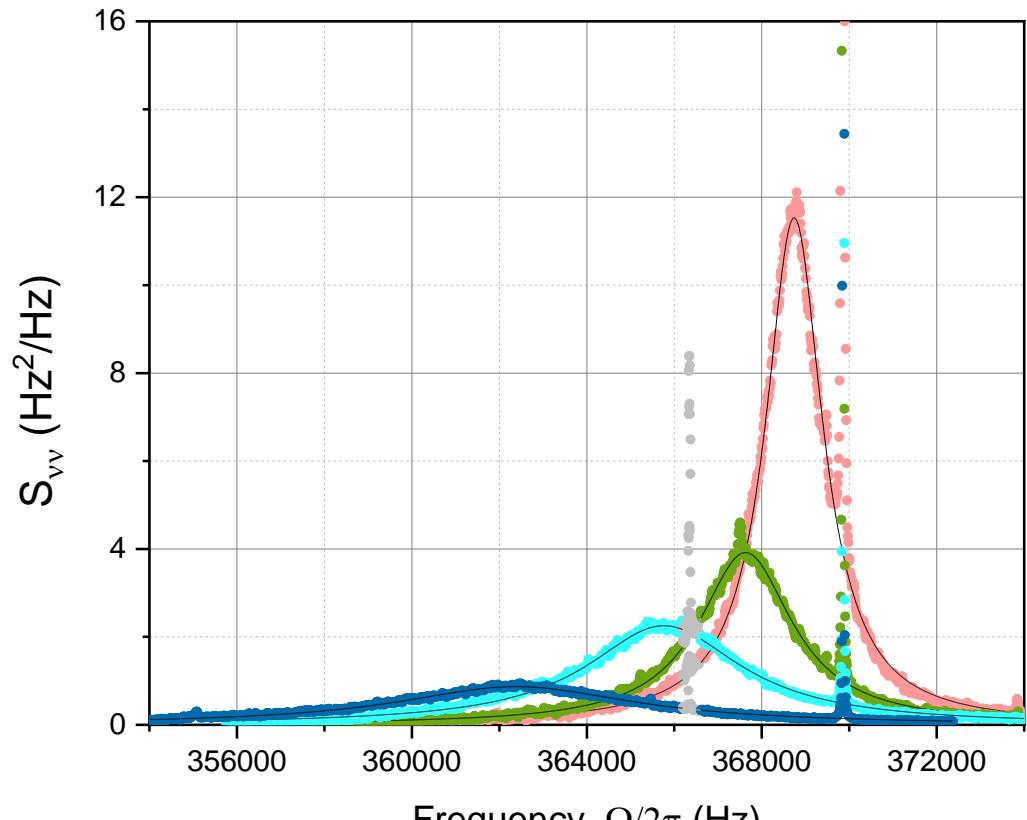
$$\Gamma_{\text{opt}}(\omega) = g^2 \frac{\Omega_m}{\omega} \left[\frac{\kappa}{(\Delta + \omega)^2 + \kappa^2/4} - \frac{\kappa}{(\Delta - \omega)^2 + \kappa^2/4} \right]$$

$$\Delta \sim \kappa/2 = 1 \text{ MHz} \quad g^2 = g_0^2 \bar{n}_{\text{cav}}$$

Cooling measurement relying on the measurement and calibration of motion induced scattering of light giving an average phonon occupancy $\langle n \rangle$

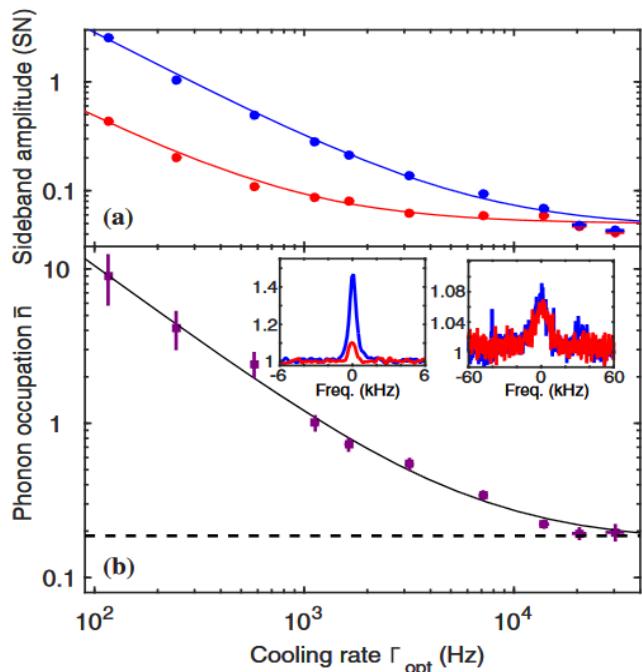
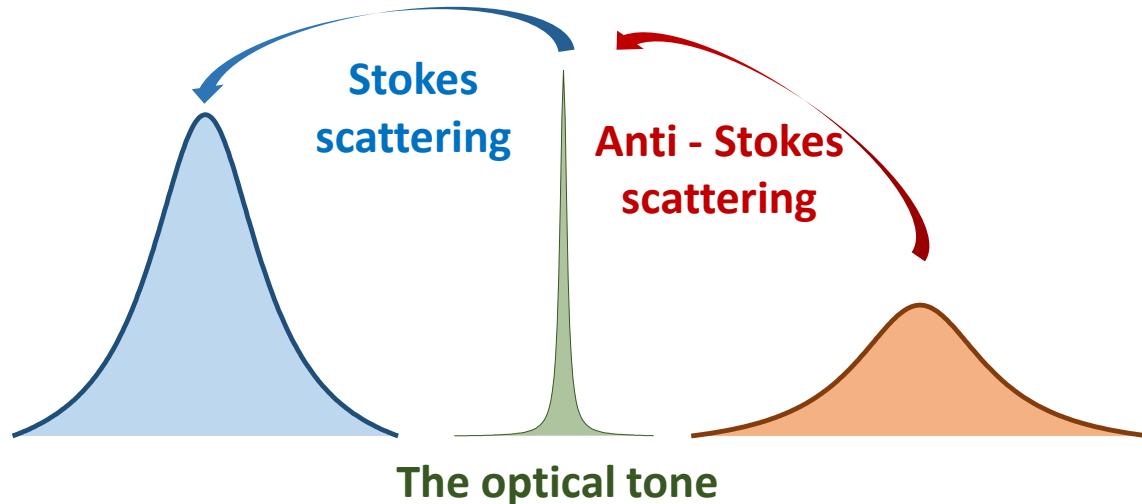
Quantum thermometry

- Ratio of Stokes and anti-Stokes sideband from the mechanical oscillator: $\frac{\langle n \rangle + 1}{\langle n \rangle}$

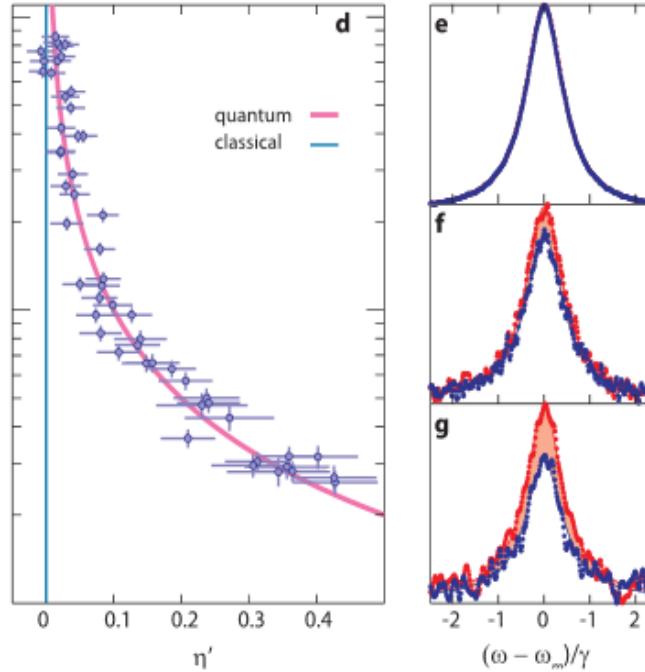


Quantum Sci. Technol. 4 (2019) 024007

The scattering picture

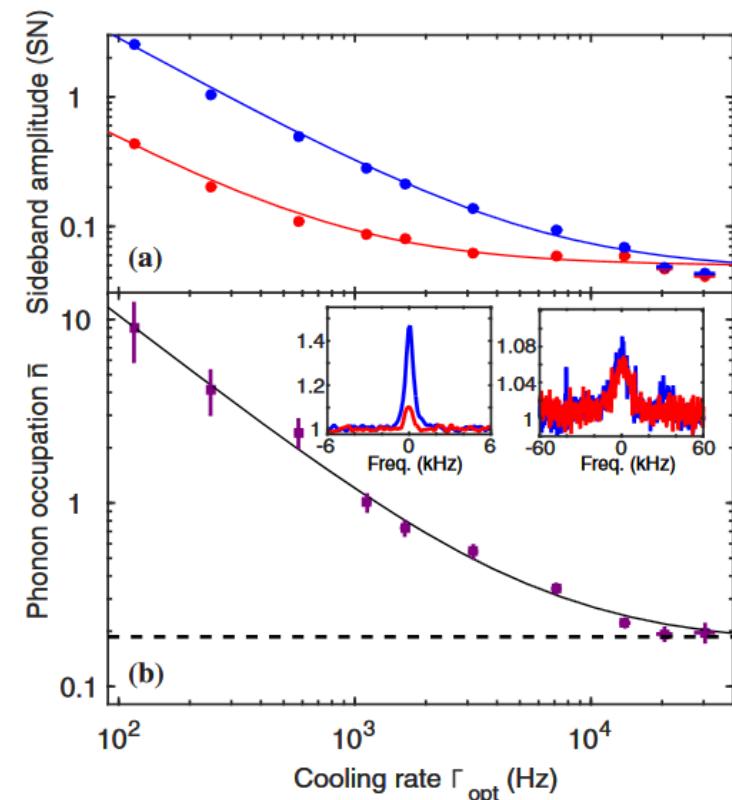


Regal, PRL 116, 063601 (2016)



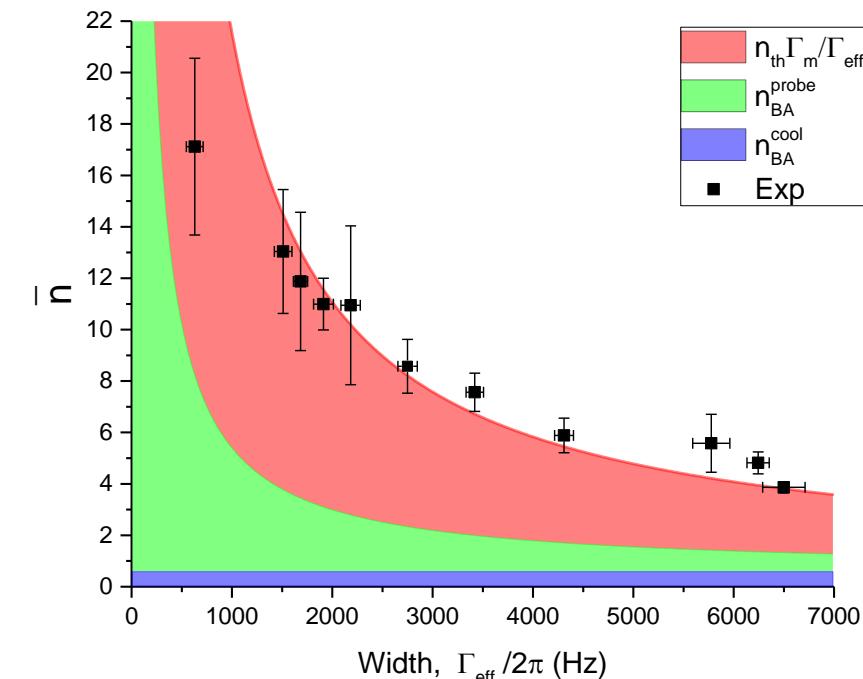
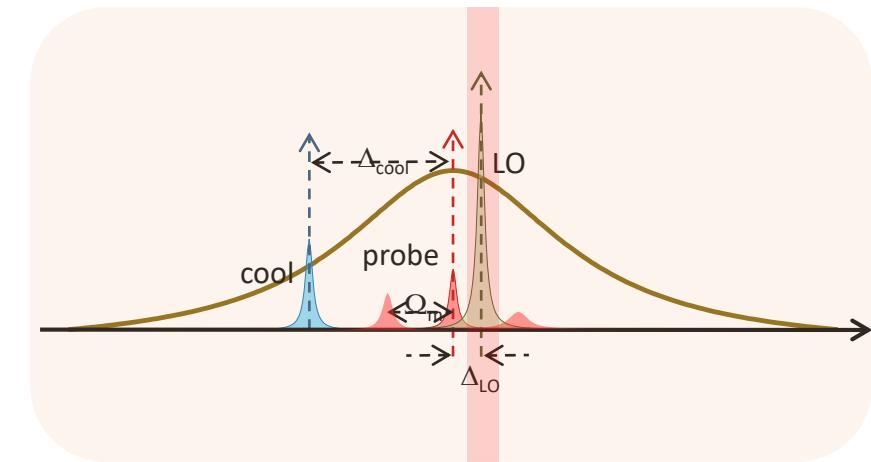
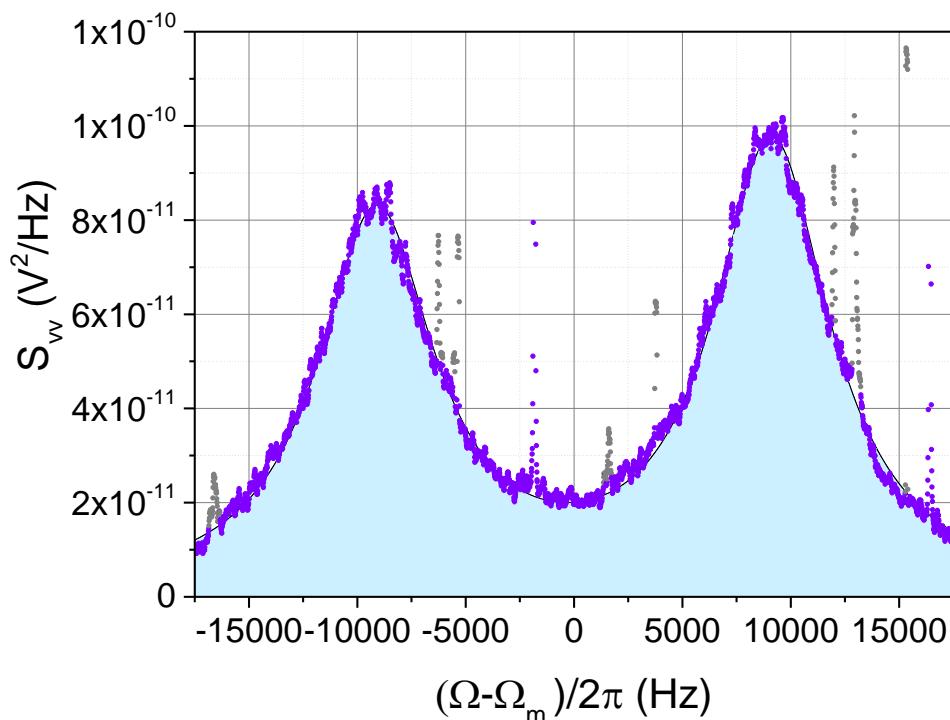
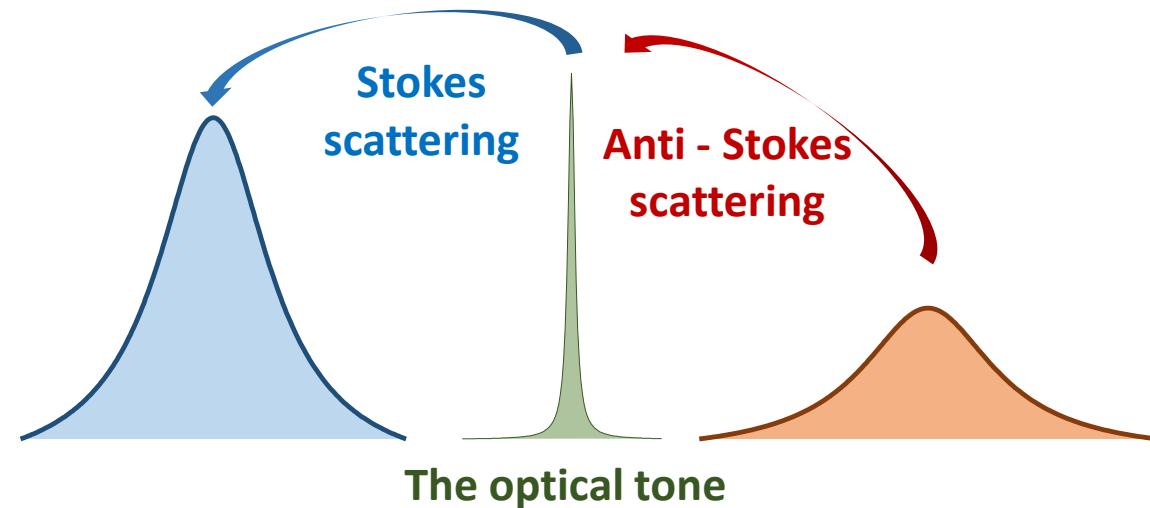
Painter, PRL 108, 033602 (2012)

$$\text{Ratio} = \frac{\text{Stokes rate}}{\text{anti Stokes rate}} = \frac{\langle n \rangle + 1}{\langle n \rangle}$$



Regal, PRL 116, 063601 (2016)

The scattering picture



Towards zero-point state

Lower bound to the position of SHO:

$$\Delta X \Delta Y \geq \frac{\hbar}{2m_{eff}\Omega_m} = x_{ZPF}$$

Typical squeezing measurement: Measurement of noise in one quadrature -> Increase of noise in the other.

Already existing squeezing schemes:

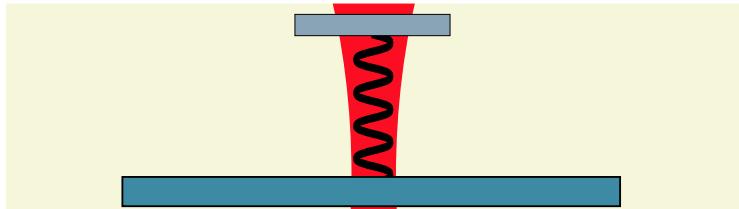
- Pulsed laser cooling of oscillator: ***requires short pulsed laser.***
- Squeezing via backaction evading measurements (BAE): ***requires motional state close to the quantum regime.***

Parametric squeezing

- No strict requirement on the starting occupation number of motional state.
- Experiments can be done in the ‘bad-cavity’ regime.

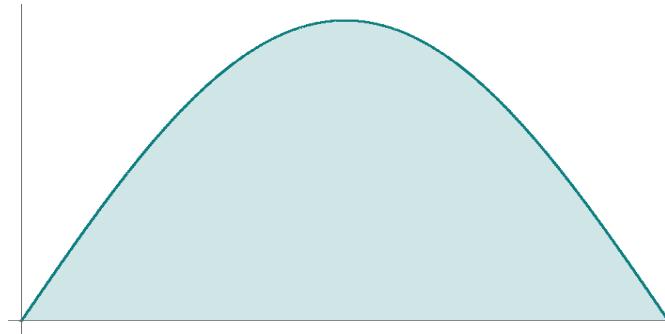
Parametric cooling of an oscillator

Modulation of spring constant at twice the resonant frequency

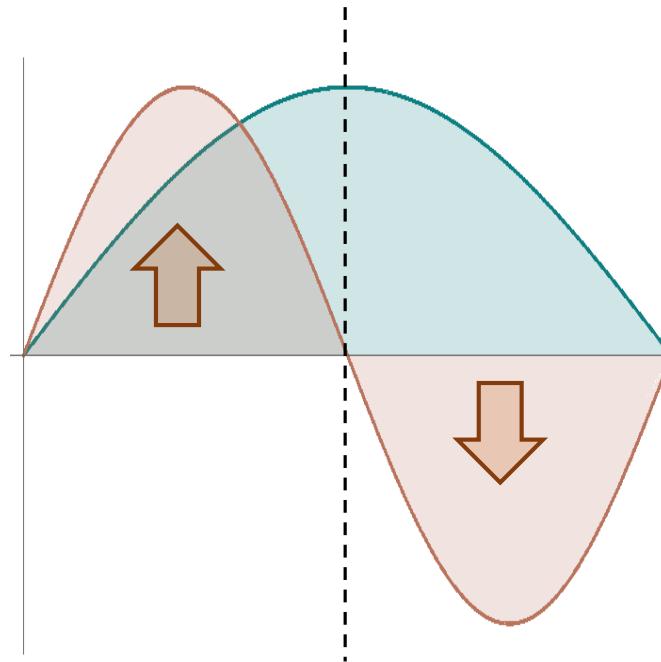


$$k(t) = k_0 + k_p(t)$$

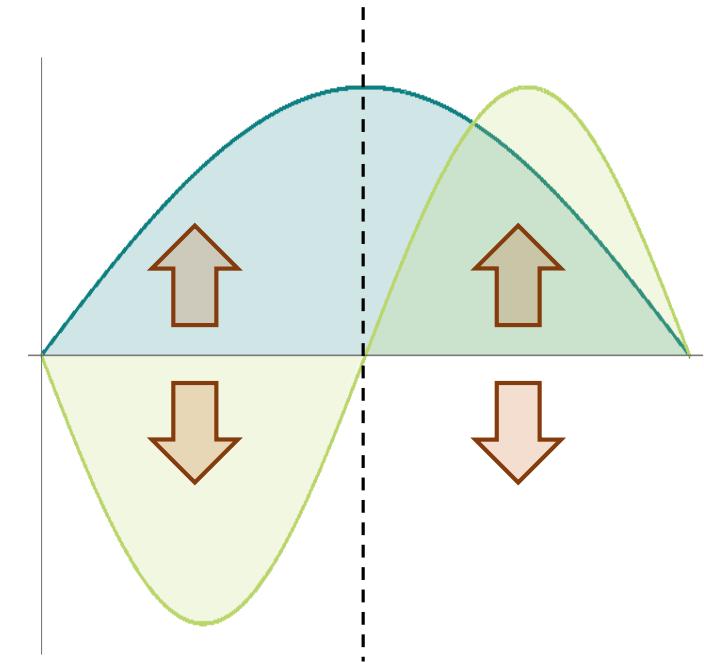
$k_0 \rightarrow$ intrinsic spring constant
 $k_p \sim \cos(2\pi f_0 t) \rightarrow$ parametric modulation



Motion of a harmonic oscillator



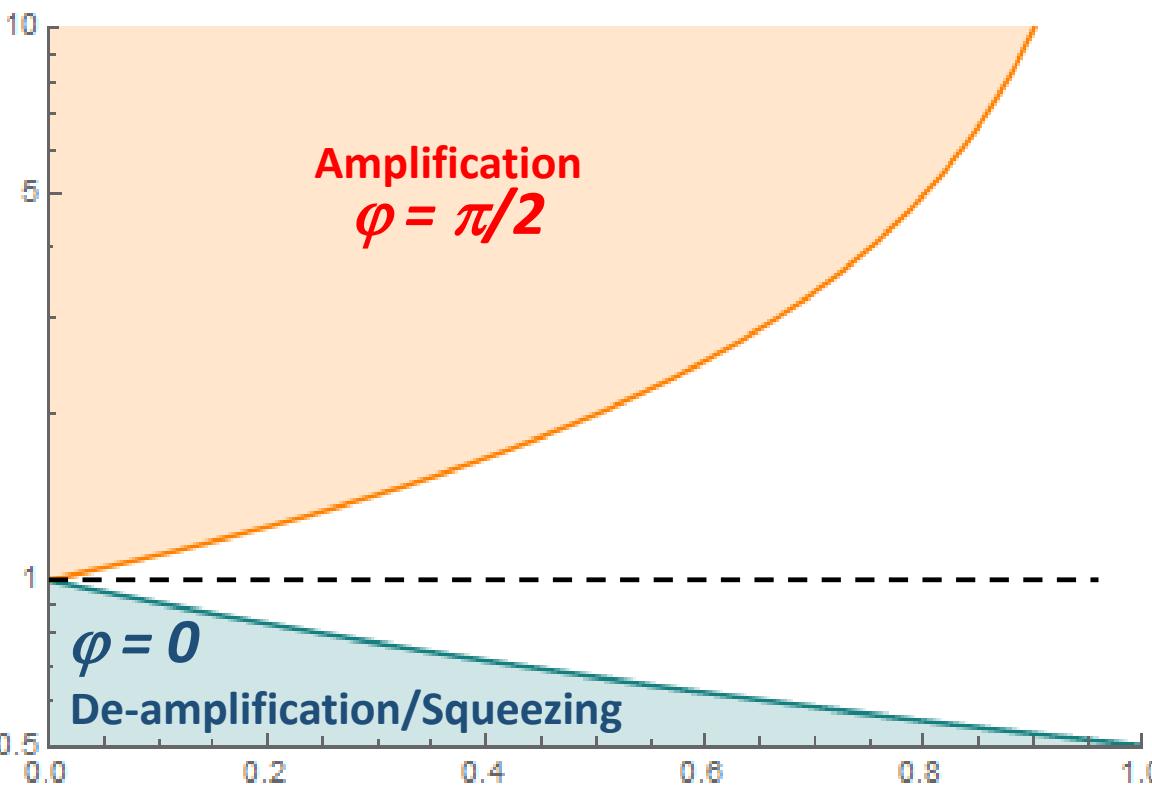
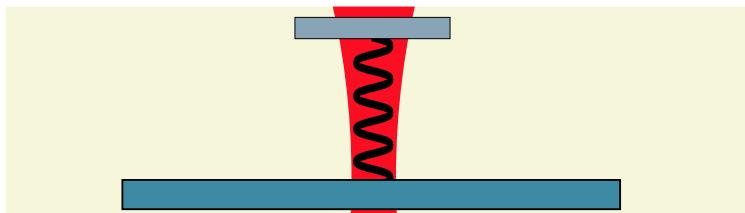
Parametric excitation @ $2f_0$



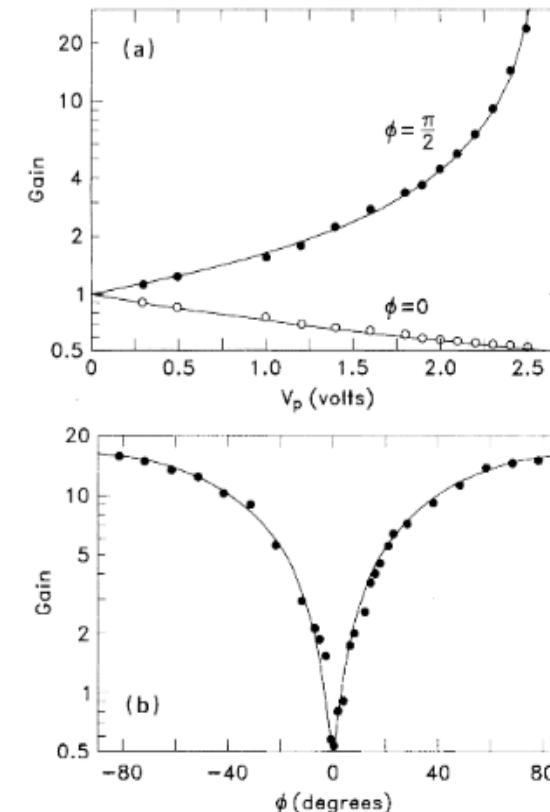
Parametric excitation @ $2f_0 + \pi$

Parametric cooling of an oscillator

Modulation of spring constant at twice the resonant frequency



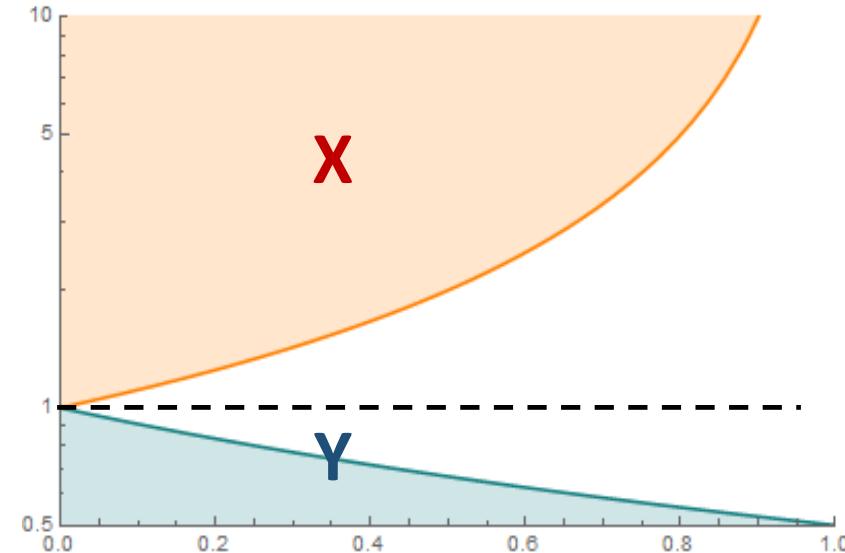
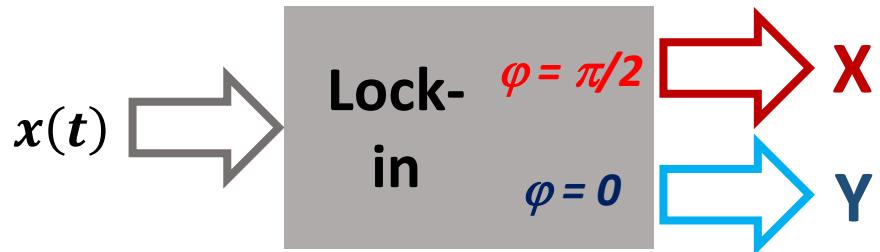
$$s(\varphi) = \left[\frac{\cos^2(\varphi)}{(1 + P_{norm})^2} + \frac{\sin^2(\varphi)}{(1 - P_{norm})^2} \right]^{1/2}$$



Rugar et al. PRL 67, 6 (1991)

Parametric cooling of an oscillator

$$x(t) = X \cos(2\pi f_0 t) + Y \cos(2\pi f_0 t)$$



Squeezing the displacement in one quadrature while adding noise on the other:
Leads to thermal squeezing of a mechanical oscillator

Parametric cooling of an oscillator

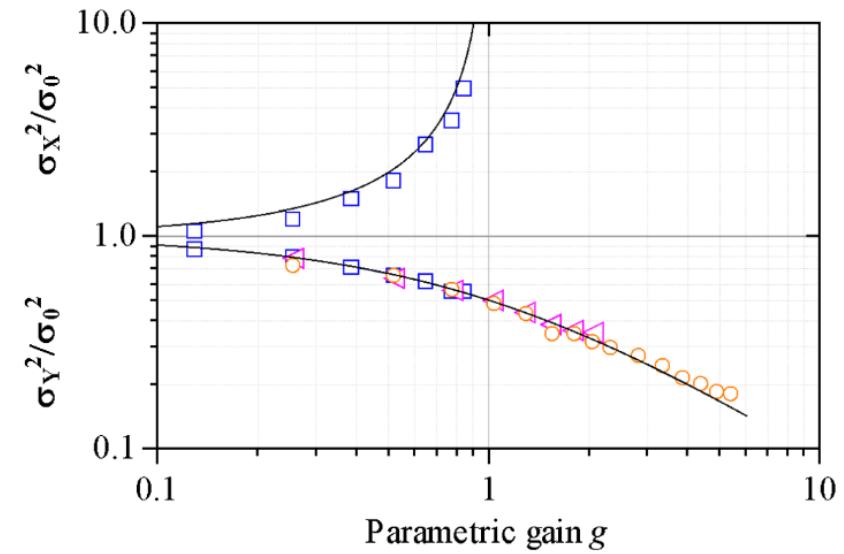
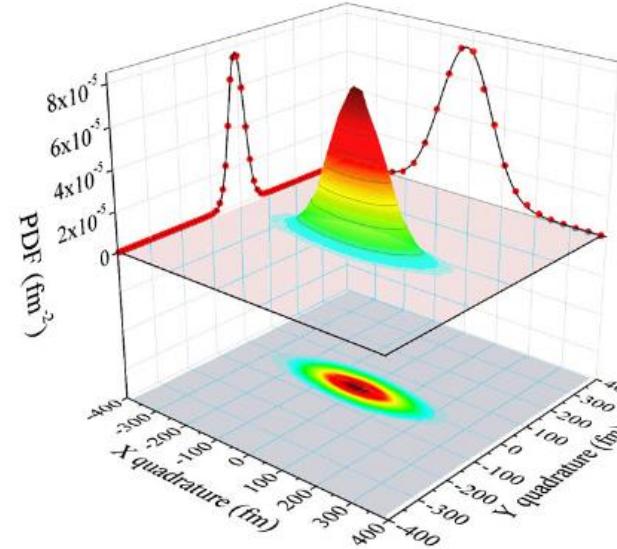
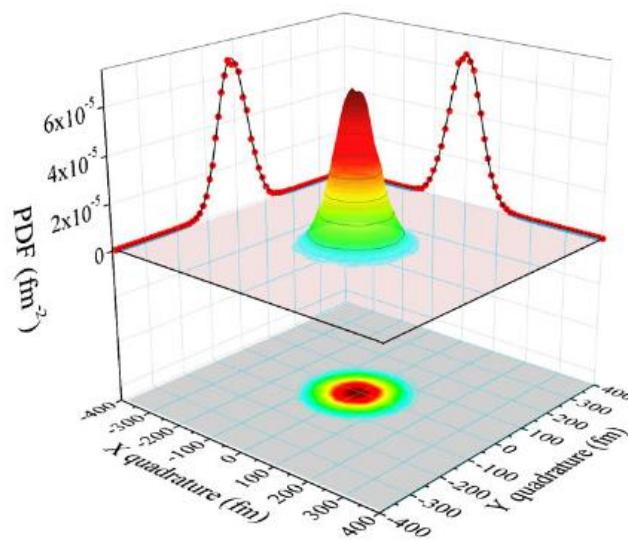
PRL 112, 023601 (2014)

PHYSICAL REVIEW LETTERS

week ending
17 JANUARY 2014

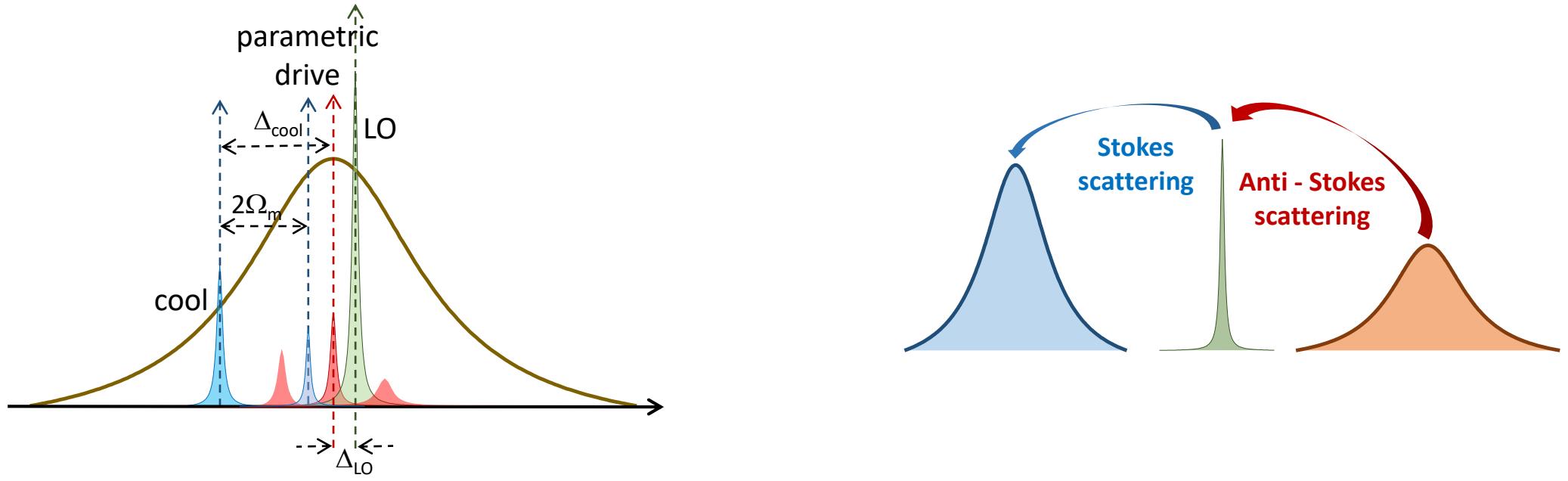
Squeezing a Thermal Mechanical Oscillator by Stabilized Parametric Effect on the Optical Spring

A. Pontin,^{1,2} M. Bonaldi,^{3,4} A. Borrielli,^{3,4} F. S. Cataliotti,^{5,6,7} F. Marino,^{7,8} G. A. Prodi,^{1,2} E. Serra,^{1,9} and F. Marin^{5,6,7,*}



Squeezed thermal state

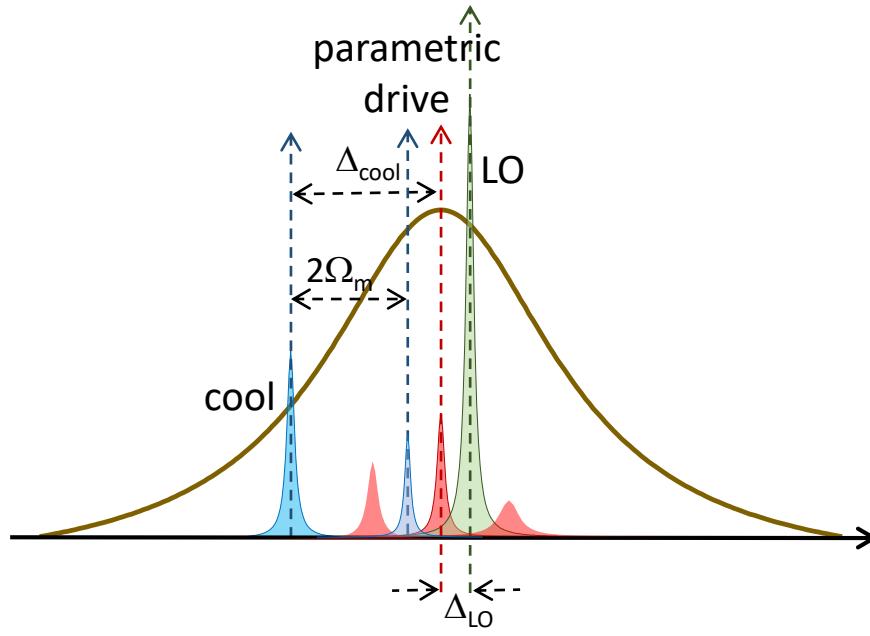
Parametric cooling close to the quantum regime



Weak parametric tone (@ $2\Omega_m$) added to the pump beam

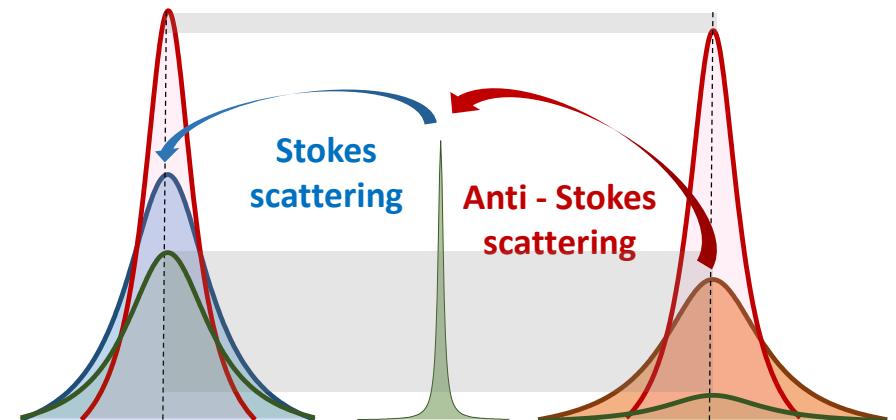
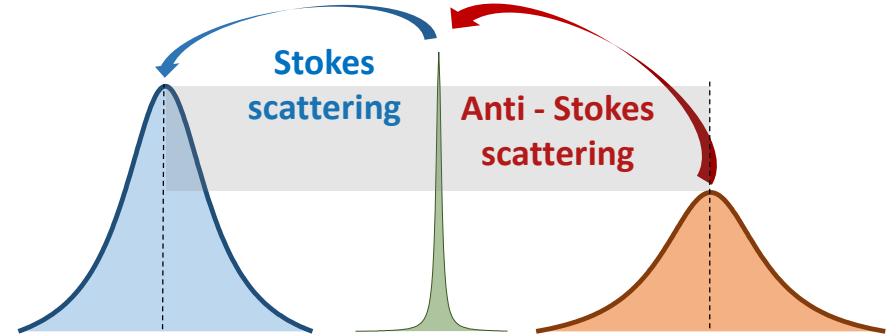
$$I_{pump} = I_{cool} + I_{par}$$

Parametric cooling close to the quantum regime

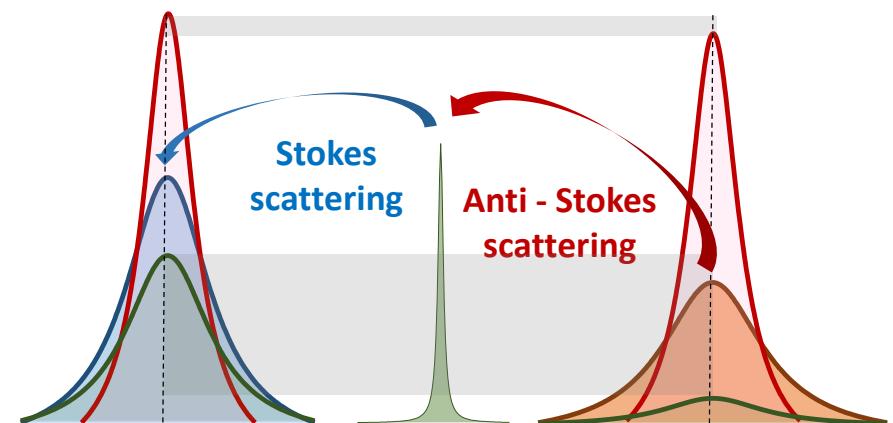
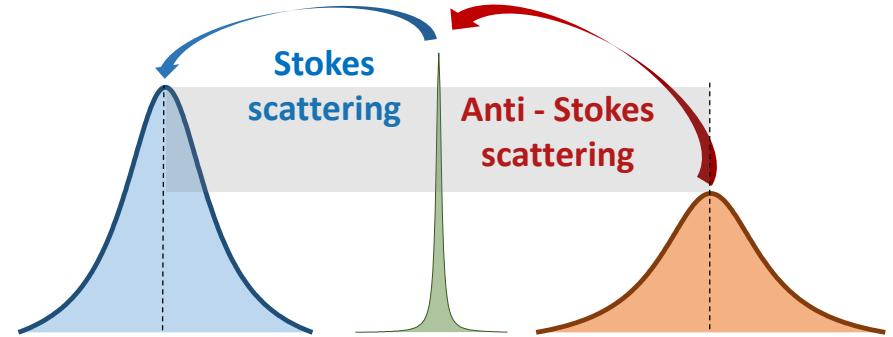
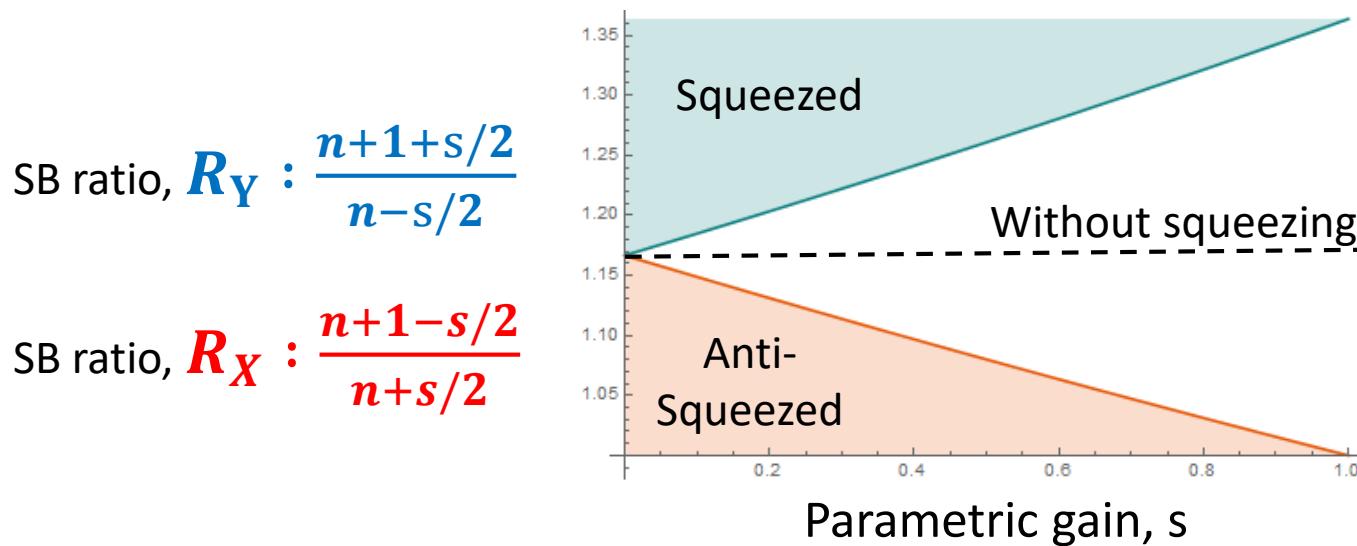
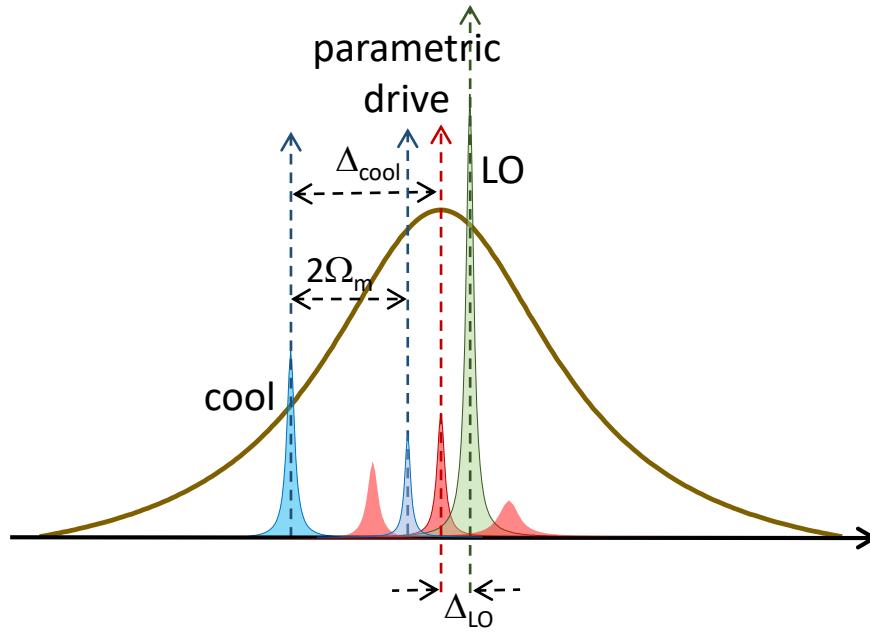


$$S_{anti-Stokes} = \frac{\Gamma_{opt}}{2} \left[\frac{n + s/2}{\omega^2 + \left(\frac{\Gamma_-}{2}\right)^2} + \frac{n - s/2}{\omega^2 + \left(\frac{\Gamma_+}{2}\right)^2} \right]$$

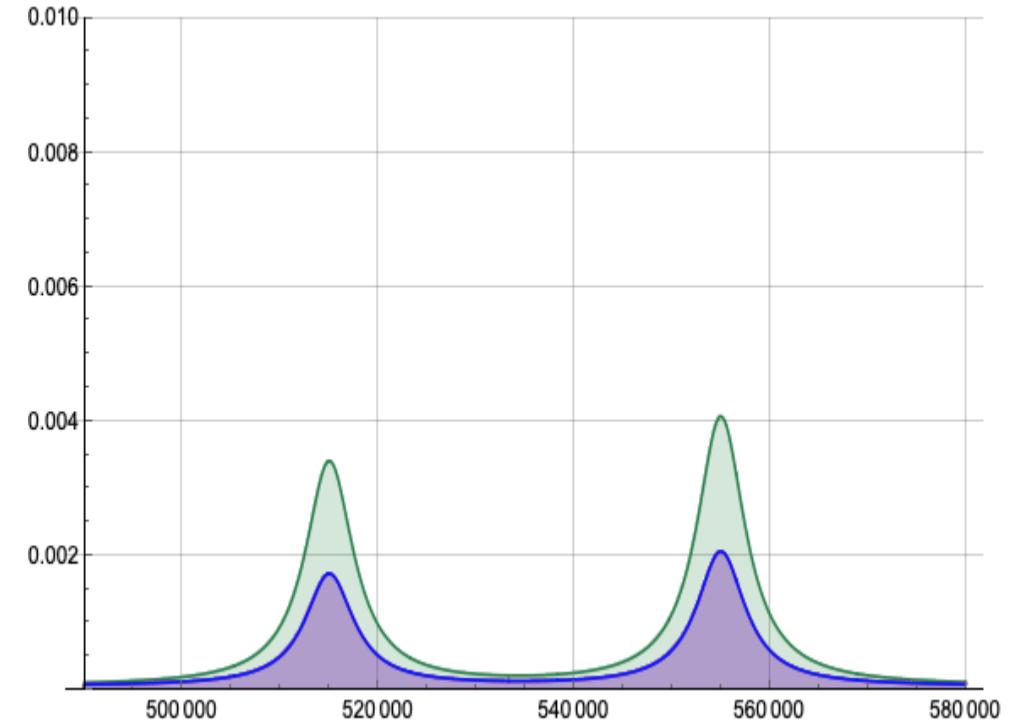
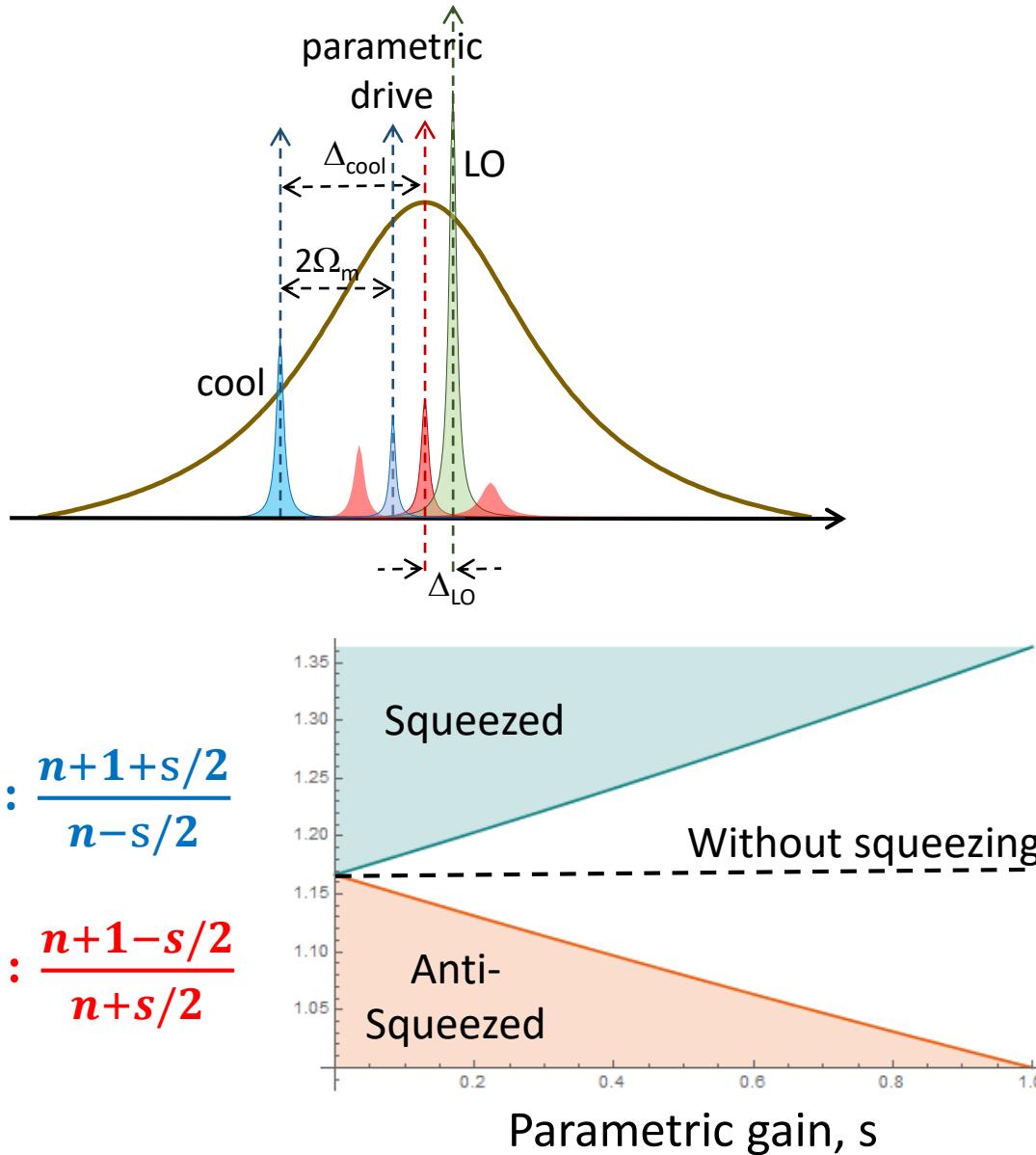
$$S_{Stokes} = \frac{\Gamma_{opt}}{2} \left[\frac{n + 1 - s/2}{\omega^2 + \left(\frac{\Gamma_-}{2}\right)^2} + \frac{n + 1 + s/2}{\omega^2 + \left(\frac{\Gamma_+}{2}\right)^2} \right]$$



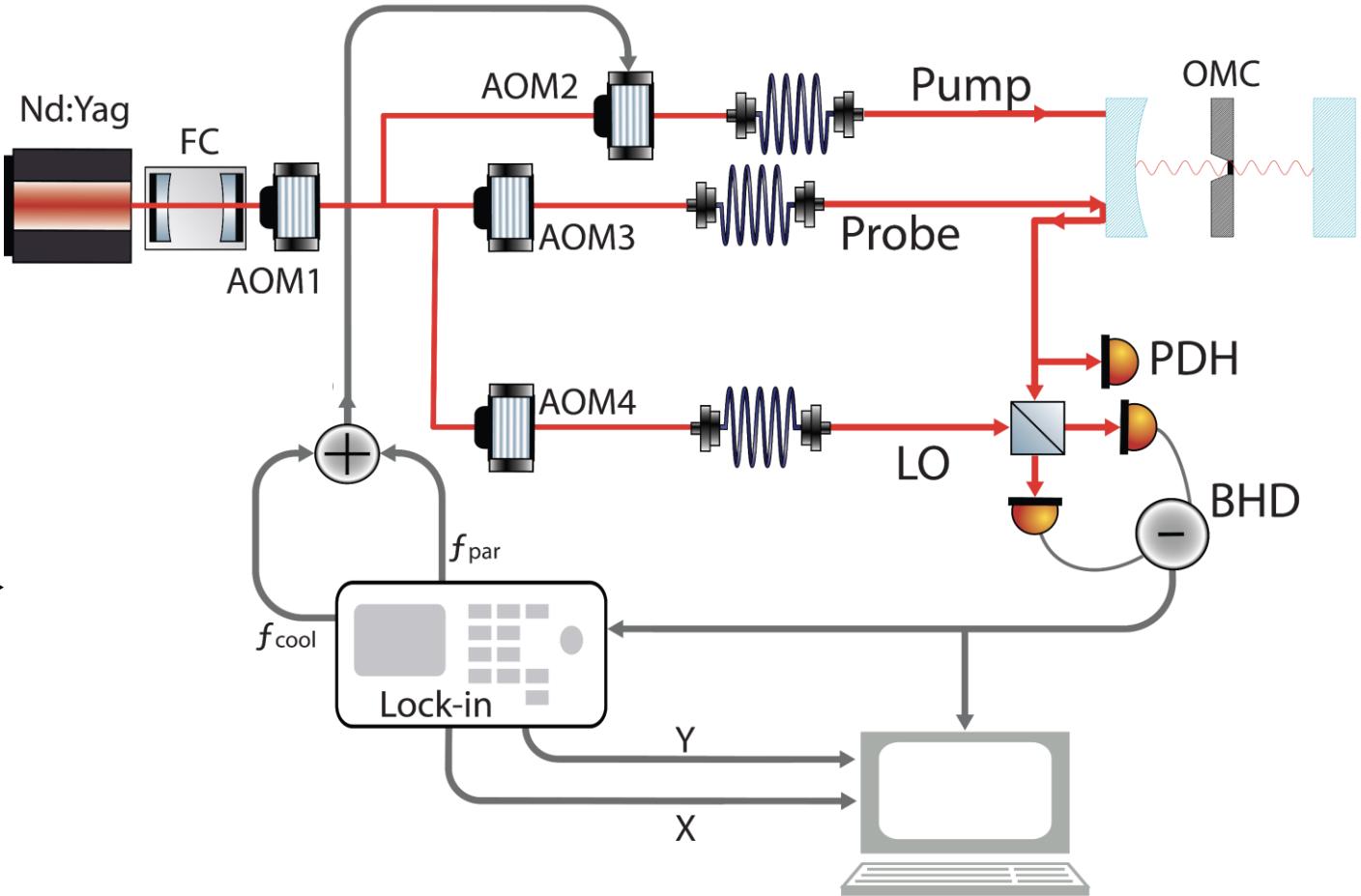
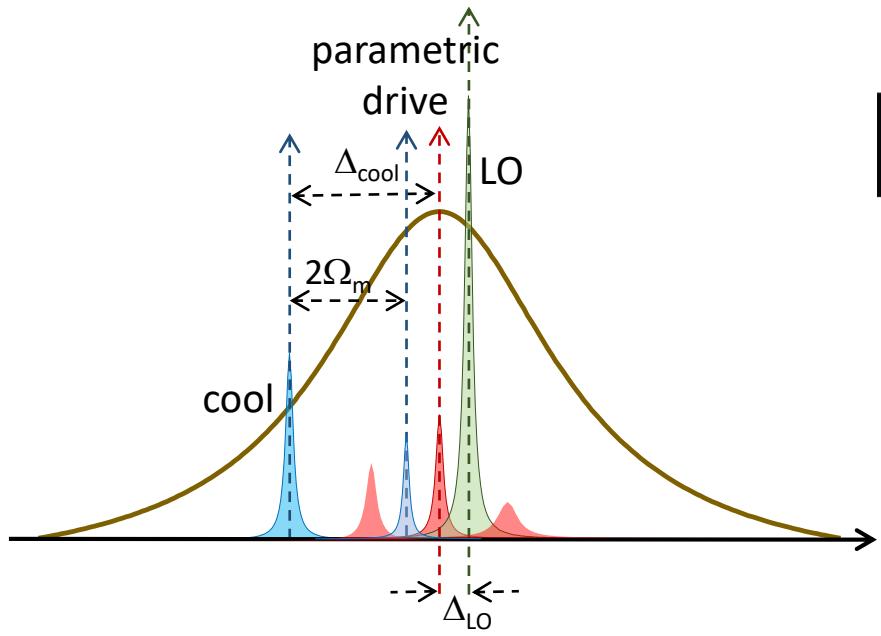
Parametric cooling close to the quantum regime



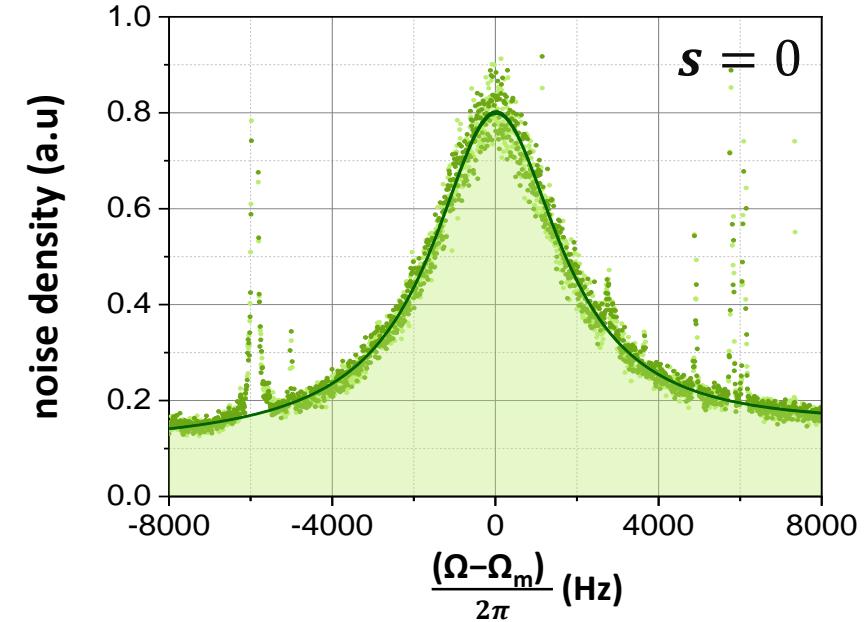
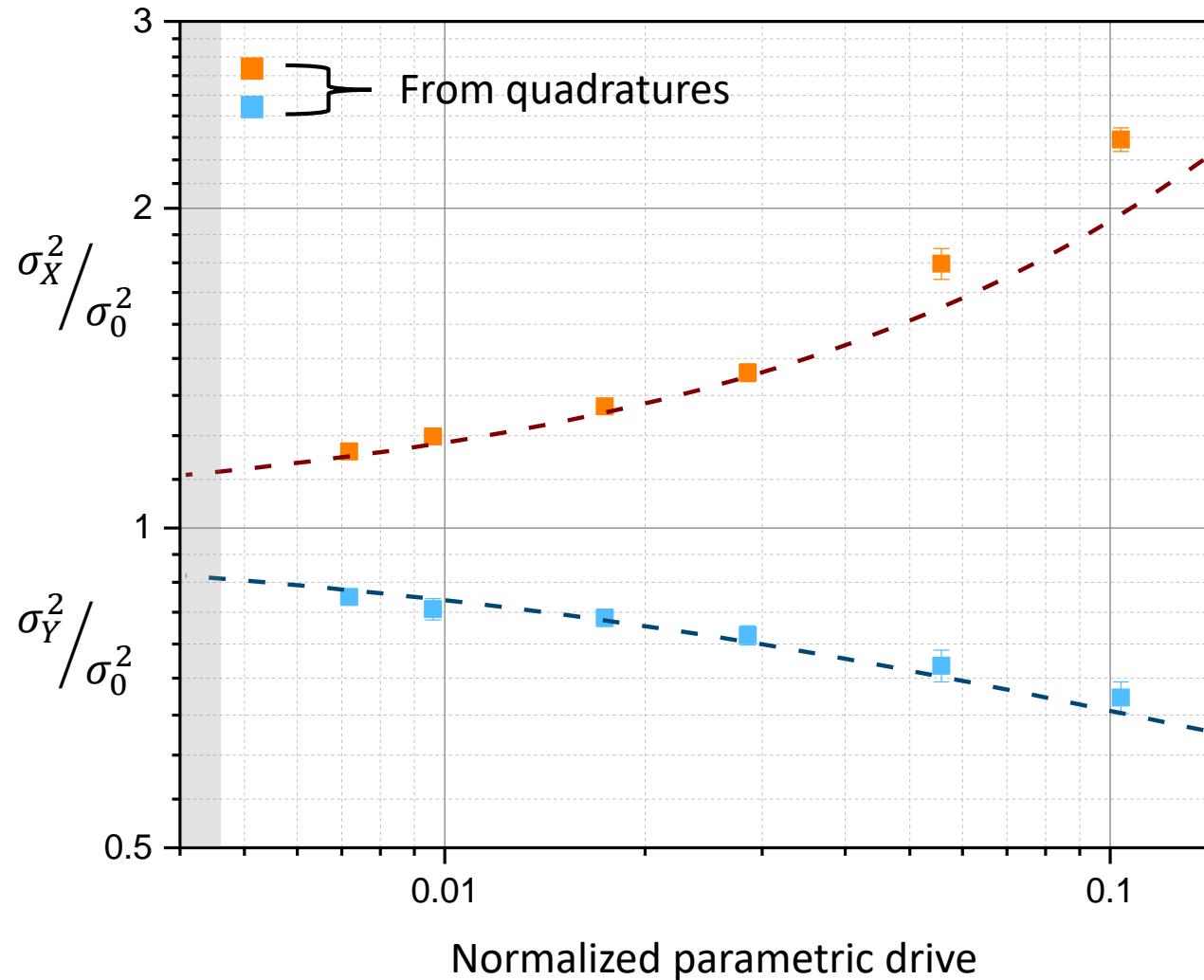
Parametric cooling close to the quantum regime



Parametric cooling: scheme

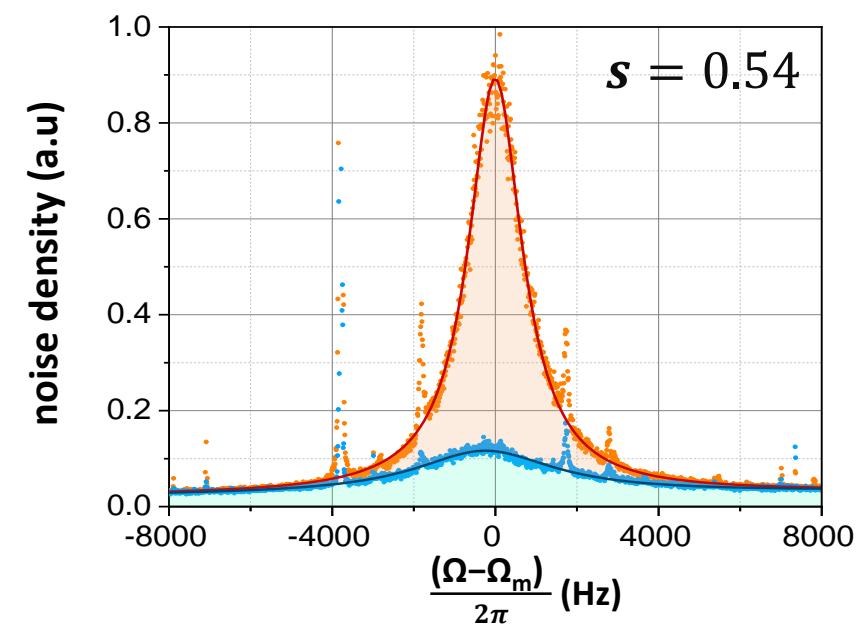
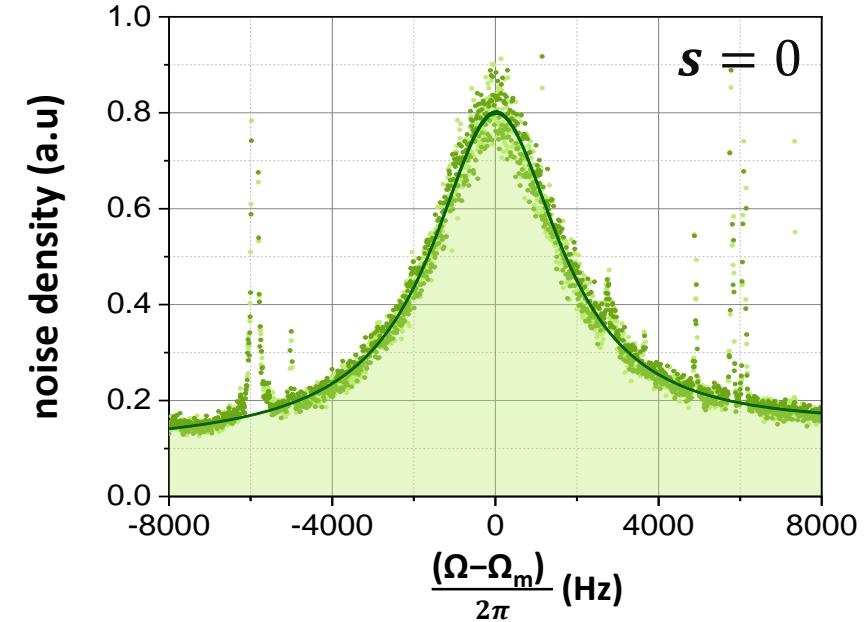
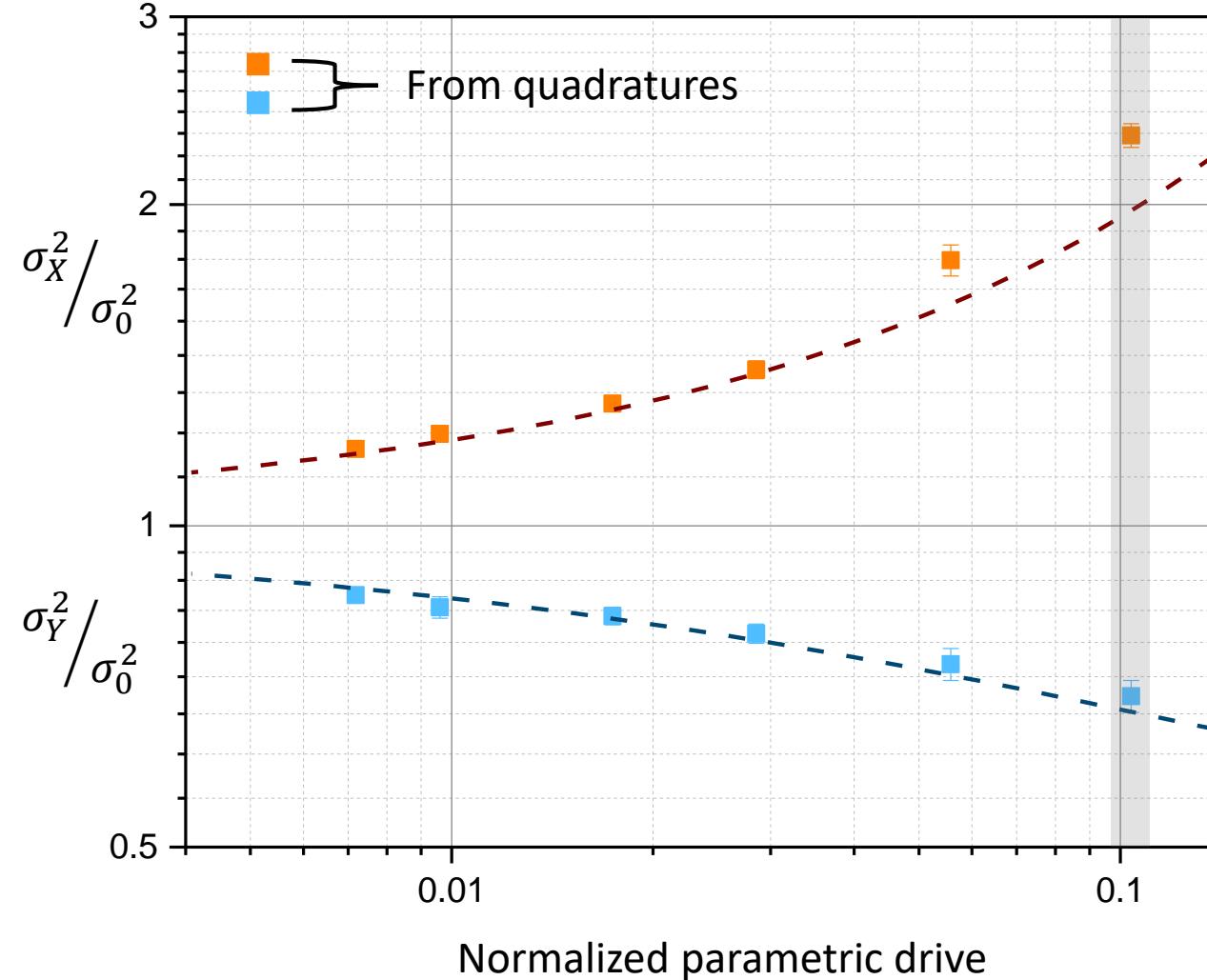


Expected variance of the quadratures



$$\sigma_X^2 = \frac{\sigma_0^2}{1-s} \quad \sigma_Y^2 = \frac{\sigma_0^2}{1+s}$$

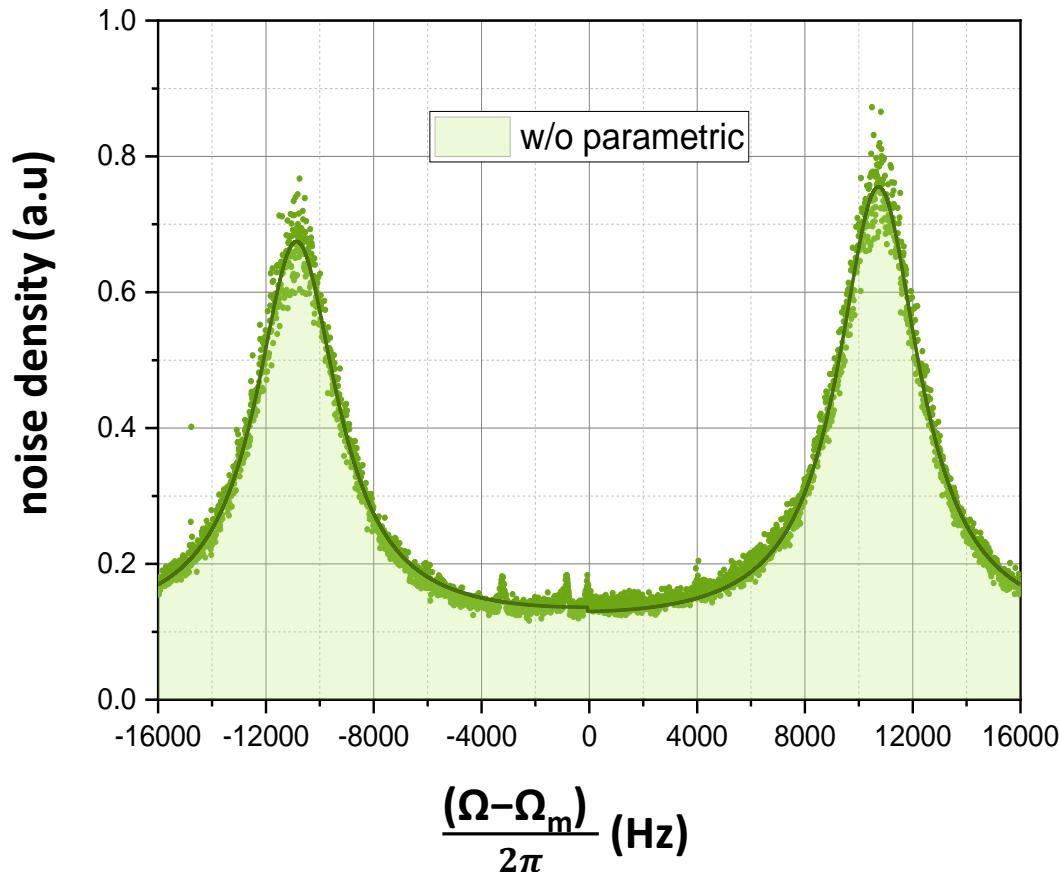
Expected variance of the quadratures



Variation of asymmetry of the quadratures

$$S_{tot} = \frac{\Gamma_{eff}}{2} \left[\frac{A_{left}}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_{eff}}{2}\right)^2} + \frac{A_{right}}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_{eff}}{2}\right)^2} \right]$$

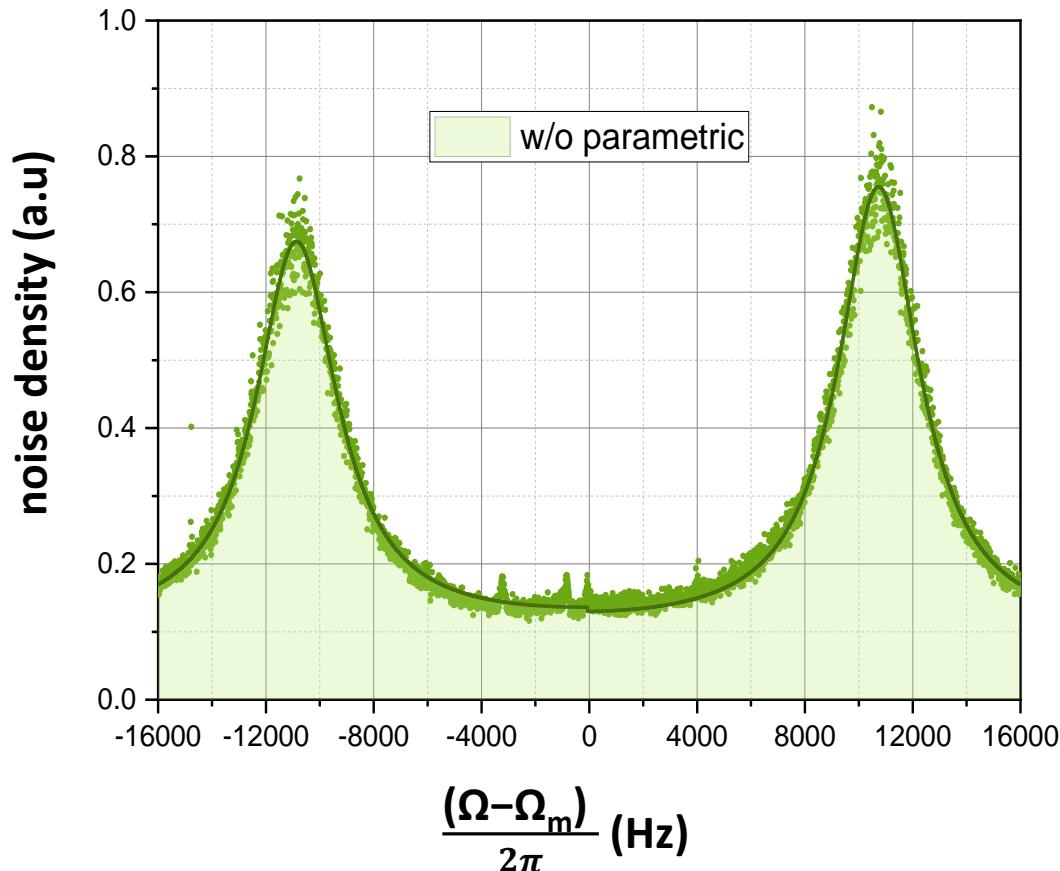
$s = 0$



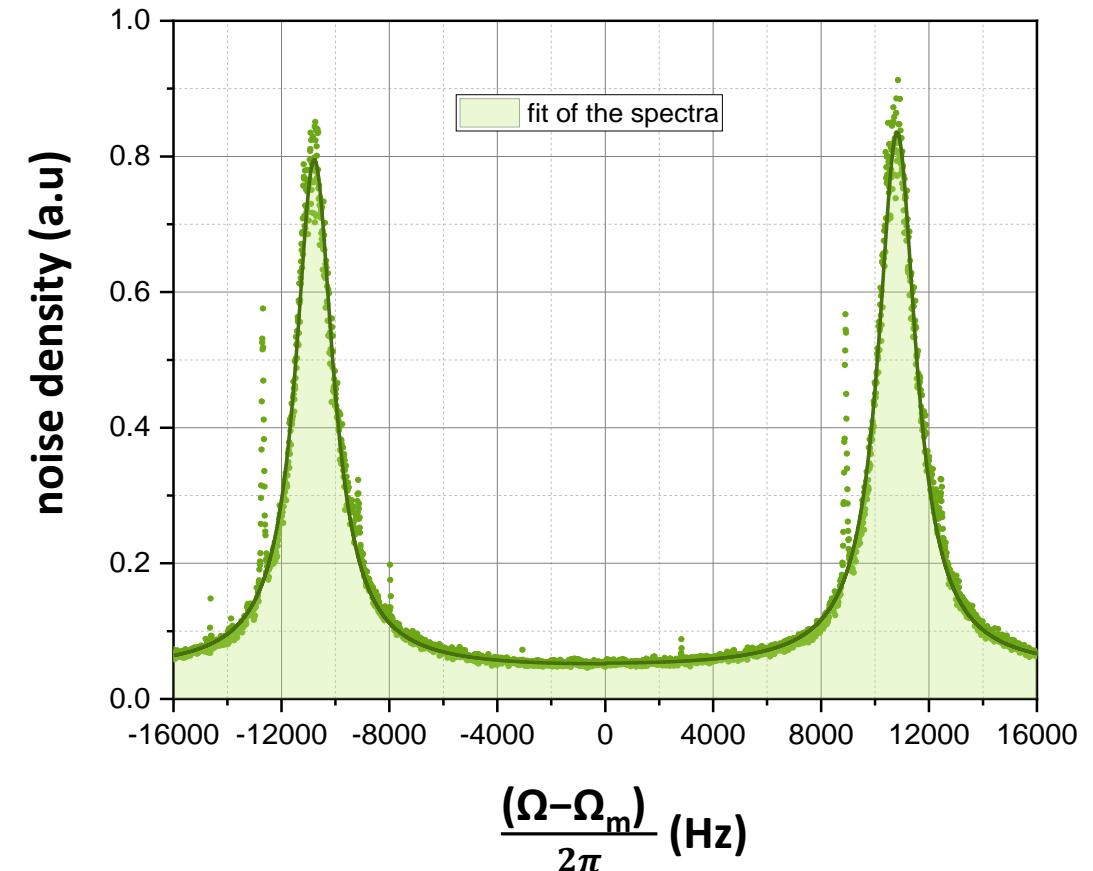
Variation of asymmetry of the quadratures

$$S_{tot} = \left[\frac{\Gamma_X A_{left}^X}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{left}^Y}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right] + \left[\frac{\Gamma_X A_{right}^X}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{right}^Y}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right]$$

$s = 0$



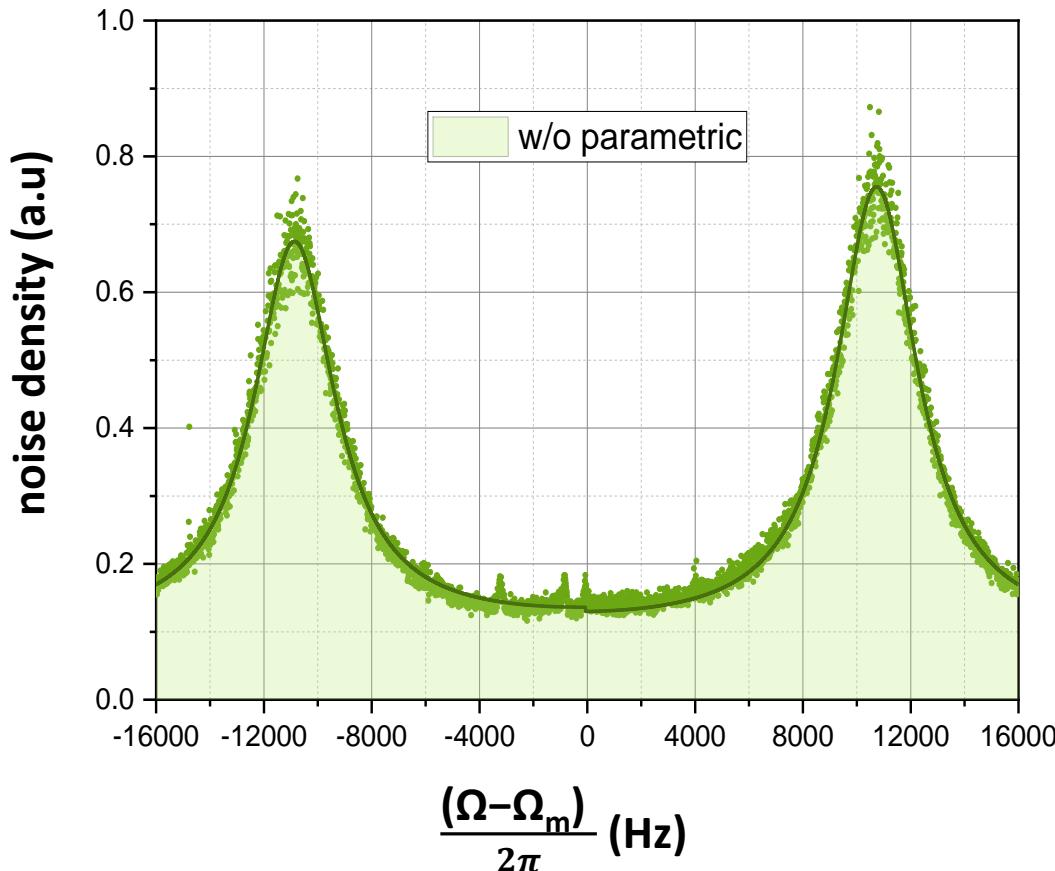
$s = 0.54$



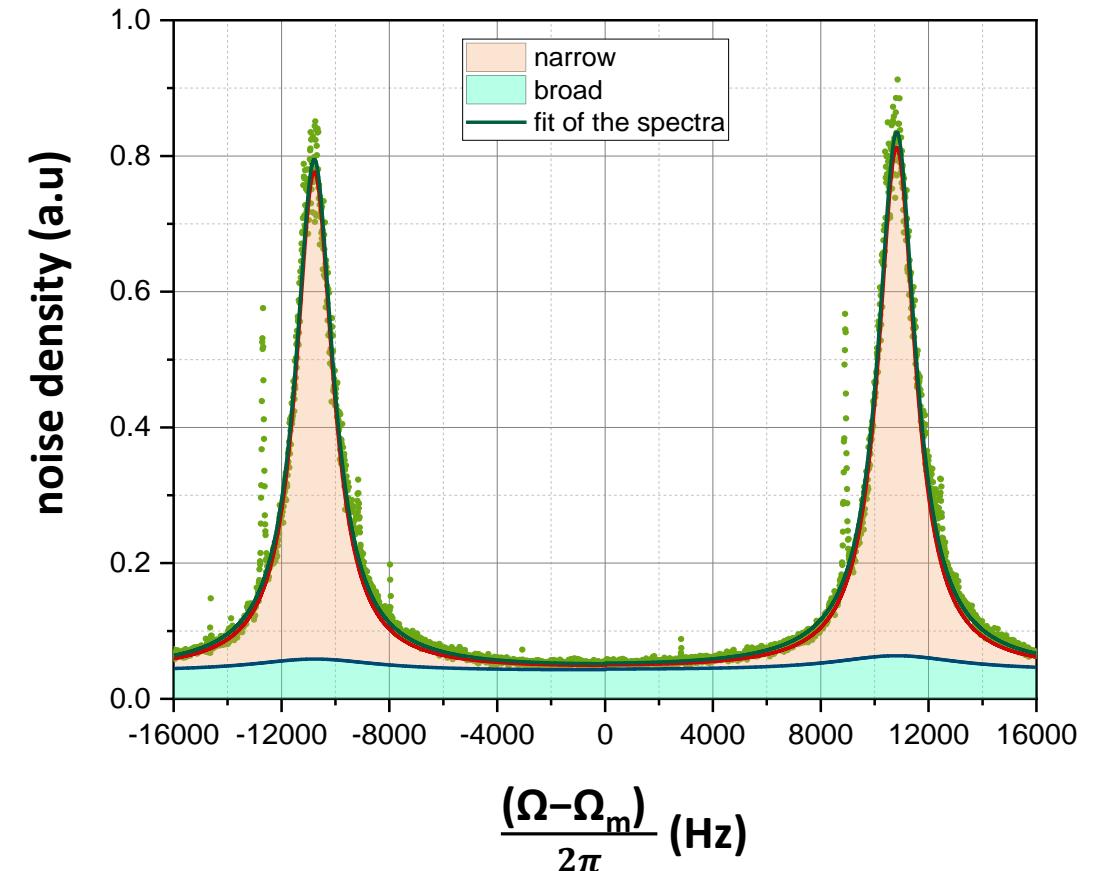
Variation of asymmetry of the quadratures

$$S_{tot} = \left[\frac{\Gamma_X A_{left}^X}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} \right] + \left[\frac{\Gamma_Y A_{left}^Y}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right] + \left[\frac{\Gamma_X A_{right}^X}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} \right] + \left[\frac{\Gamma_Y A_{right}^Y}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right]$$

$s = 0$



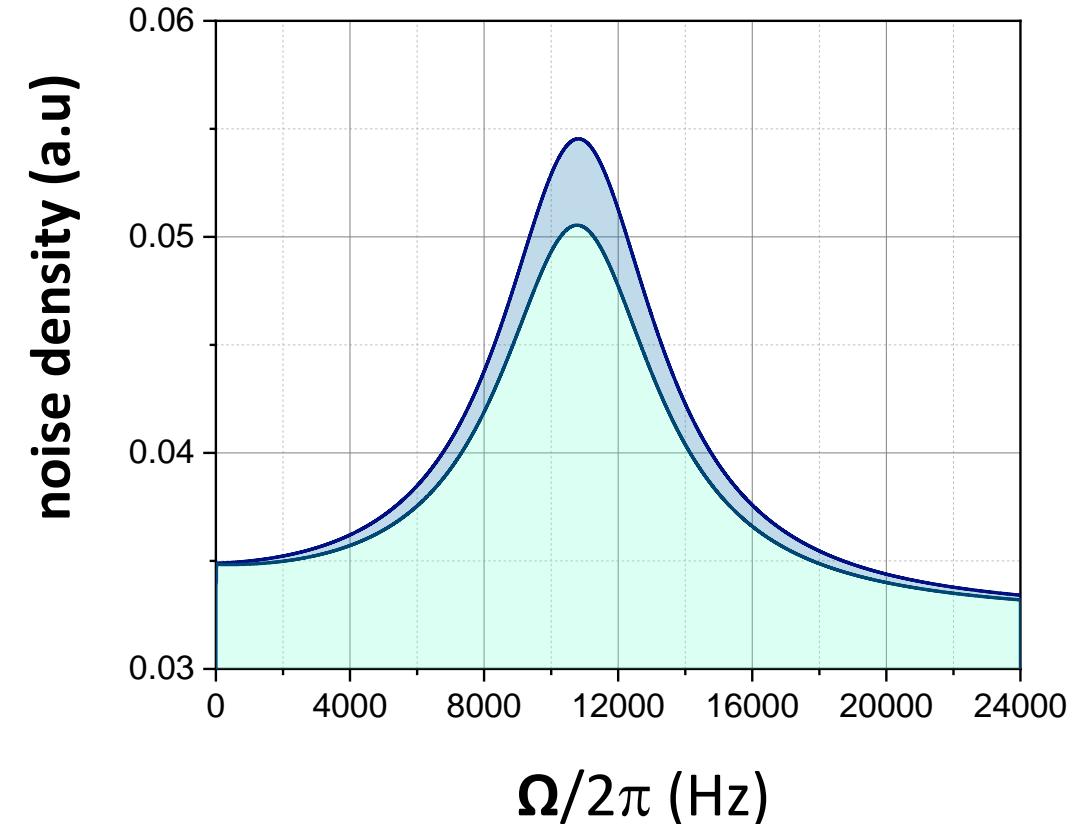
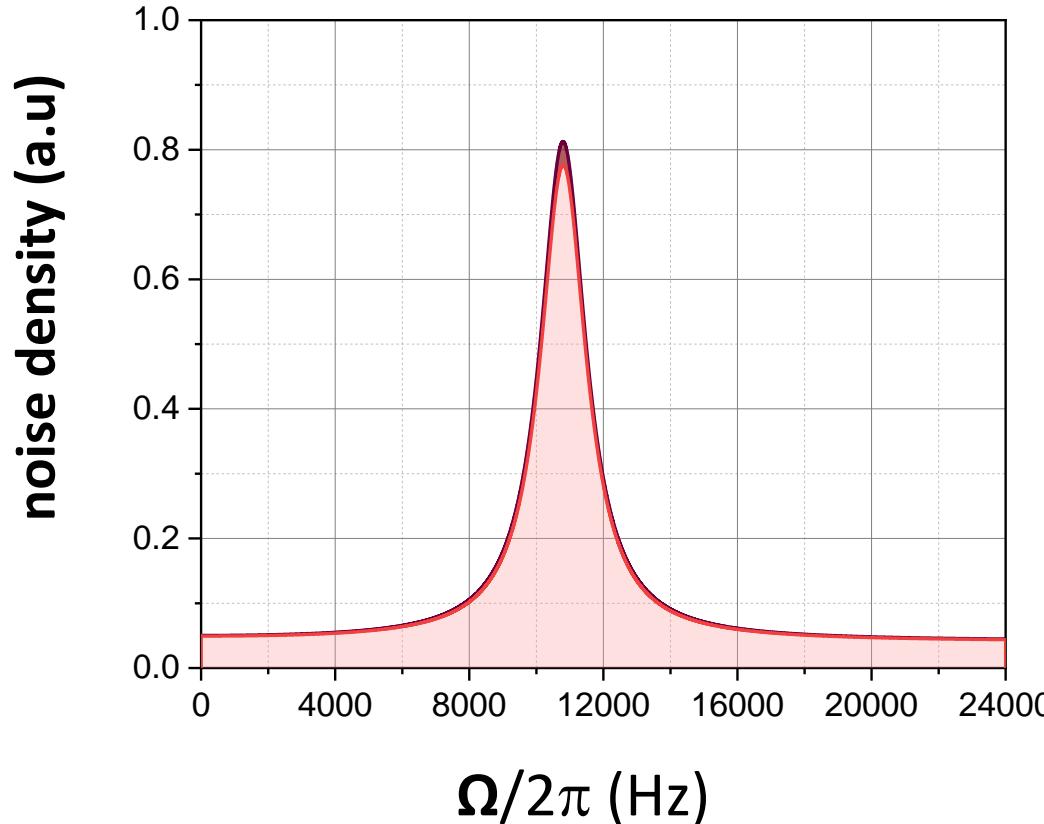
$s = 0.54$



Variation of asymmetry of the quadratures

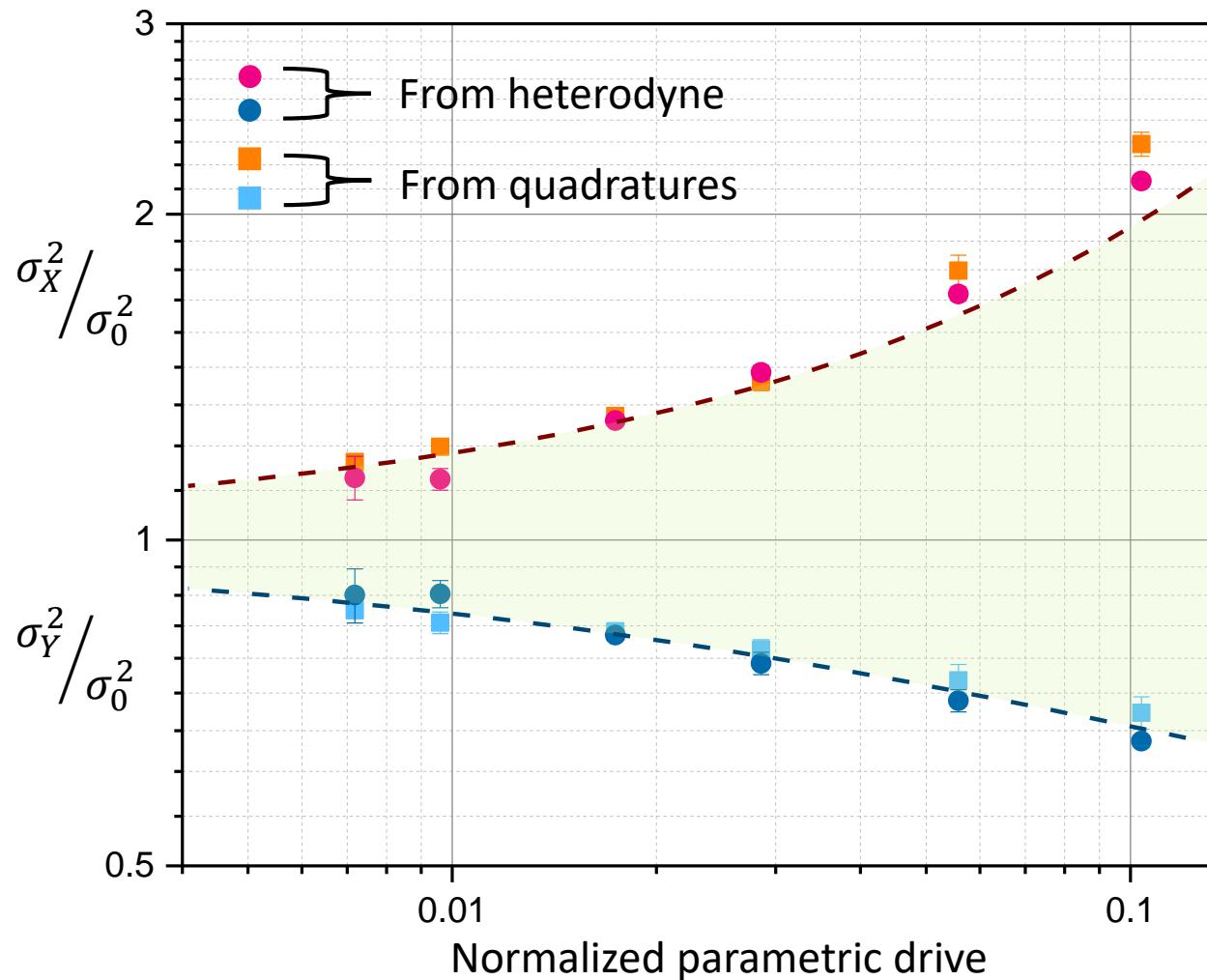
$$S_{tot} = \left[\frac{\Gamma_X A_{left}^X}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{left}^Y}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right] + \left[\frac{\Gamma_X A_{right}^X}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{right}^Y}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right]$$

$$s = 0.54$$



Variation of asymmetry of the quadratures

$$S_{tot} = \left[\frac{\Gamma_X A_{left}^X}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{left}^Y}{(\omega - \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} + \frac{\Gamma_X A_{right}^X}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_X}{2}\right)^2} + \frac{\Gamma_Y A_{right}^Y}{(\omega + \Omega_{LO})^2 + \left(\frac{\Gamma_Y}{2}\right)^2} \right]$$



$$\Gamma_X = \Gamma_{eff}(1 - s)$$

$$\Gamma_Y = \Gamma_{eff}(1 + s)$$

$$\sigma_X^2 = \frac{\sigma_0^2}{1 - s}$$

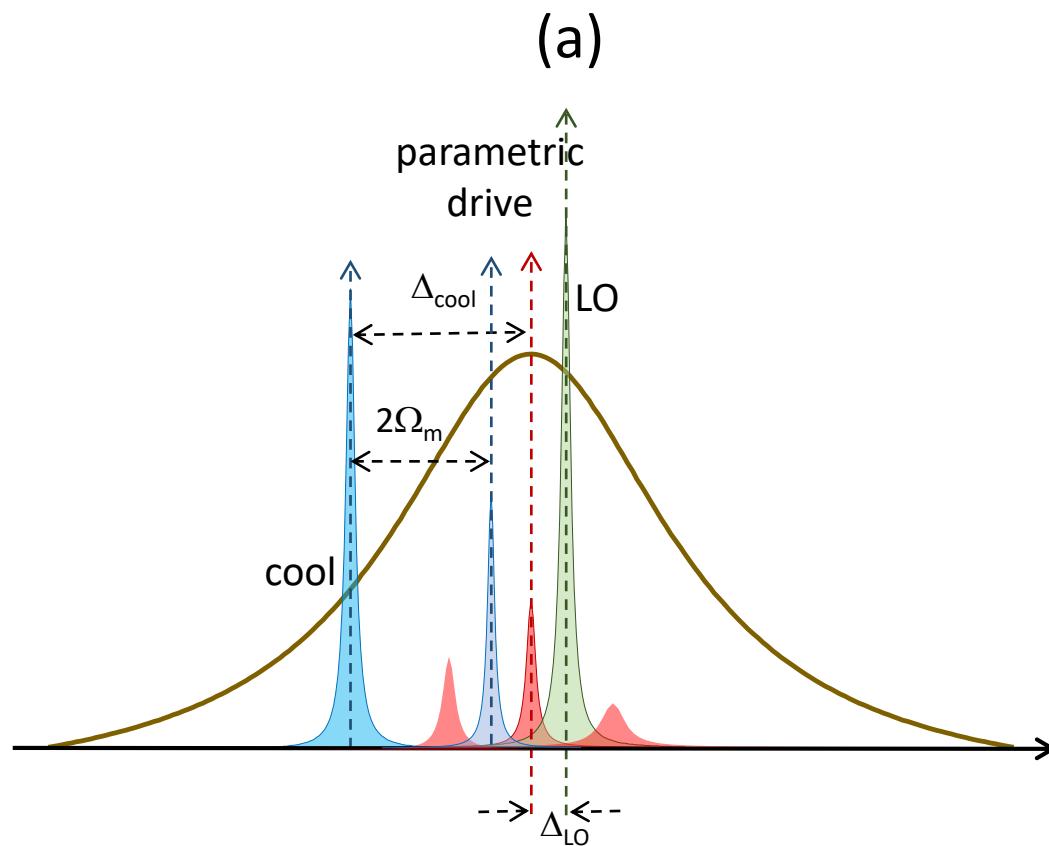
$$\sigma_Y^2 = \frac{\sigma_0^2}{1 + s}$$

Variation of asymmetry of the quadratures

$$I_{pump} = I_{cool} + I_{par}, \text{ where: } I_{par} = \alpha I_{pump}$$

(a) $I_{pump} \uparrow$ keeping ' α ' constant: 's' is constant

(b) $I_{par} \uparrow$ keeping I_{pump} constant: 's' varies keeping Γ_{eff} constant

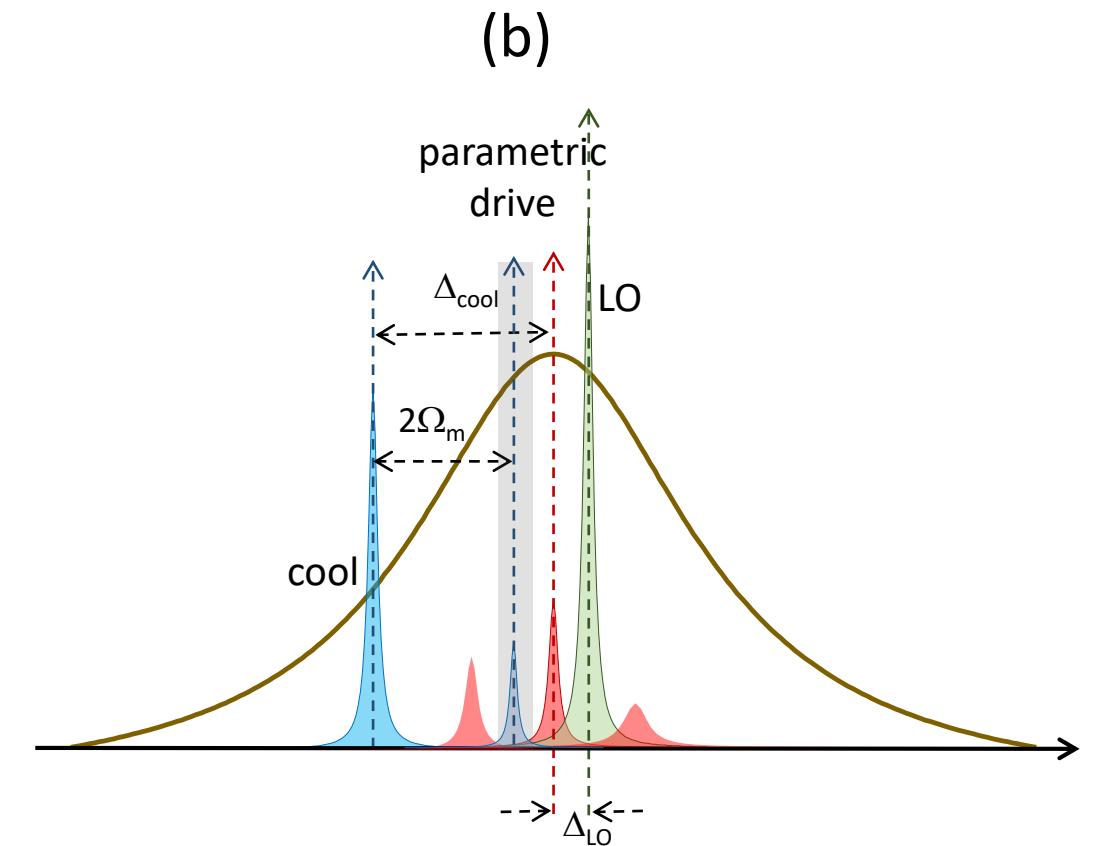
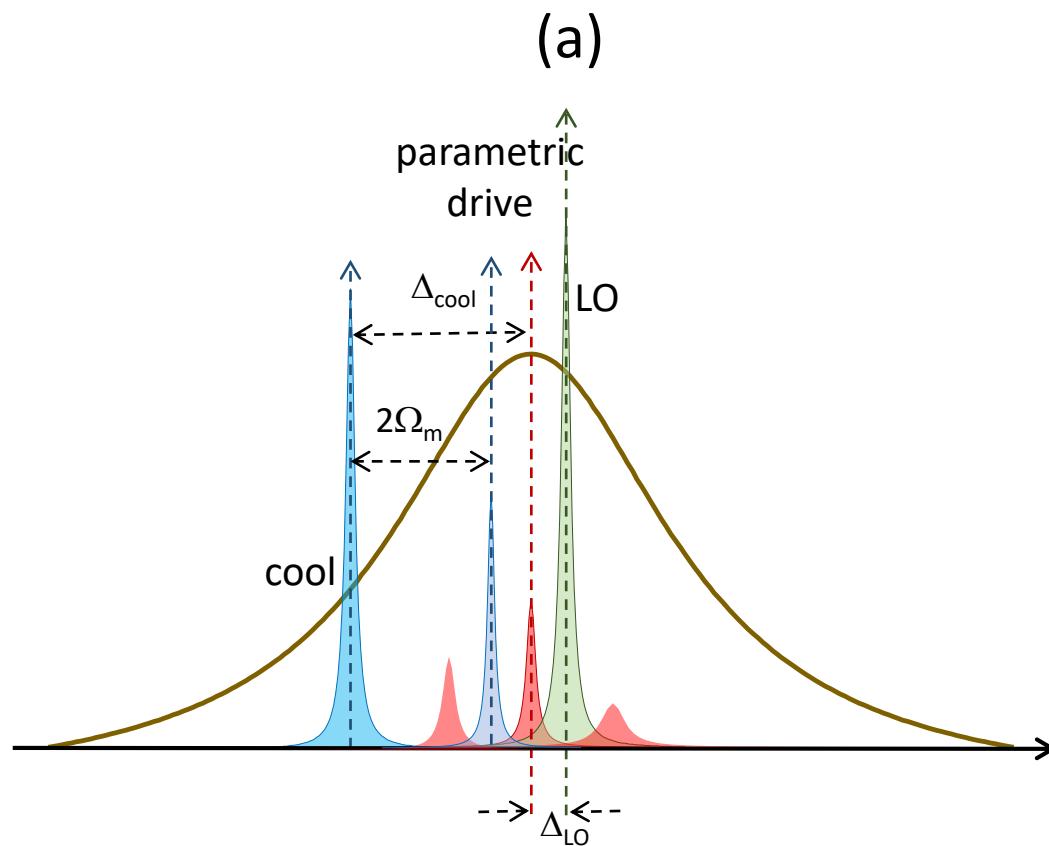


Variation of asymmetry of the quadratures

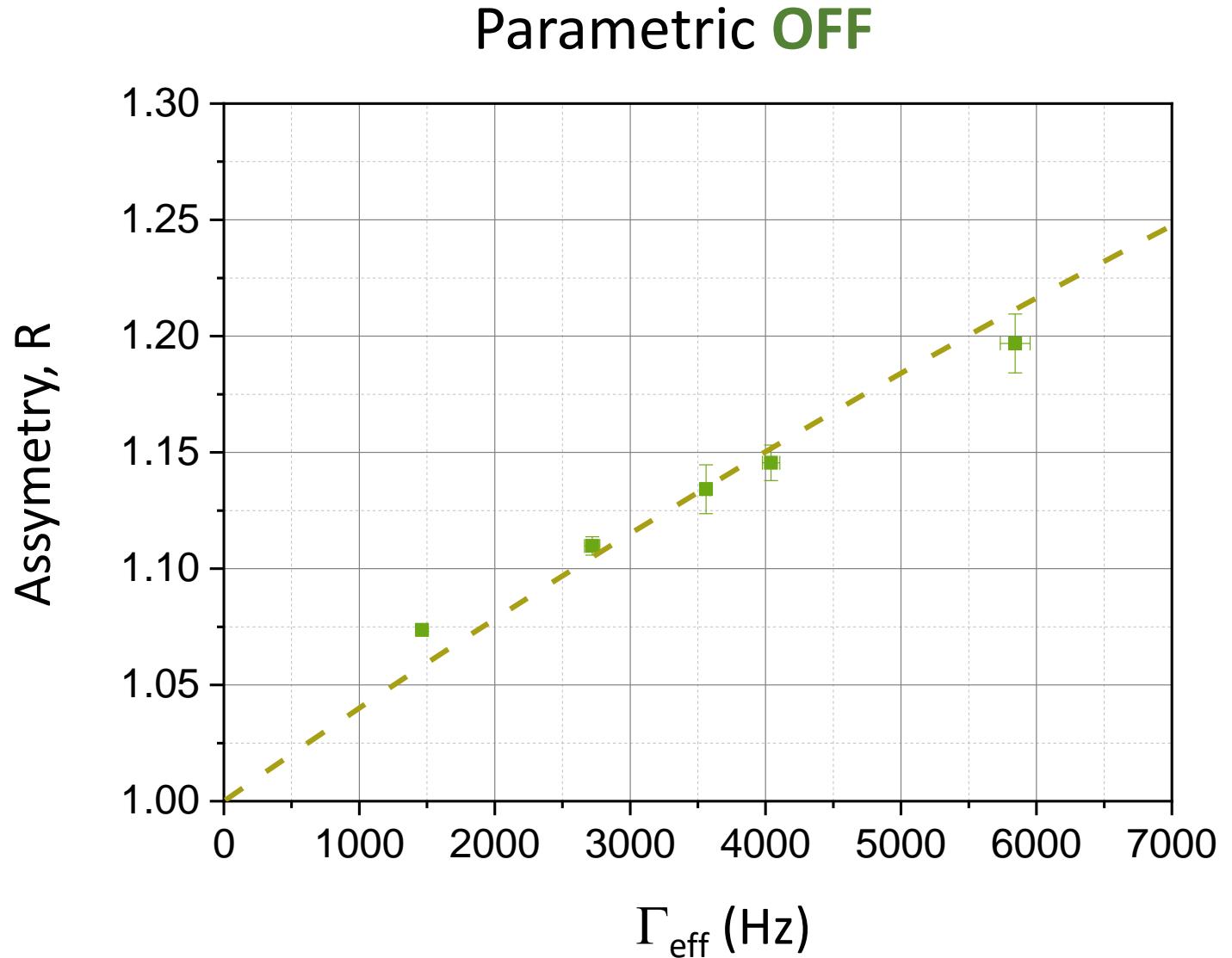
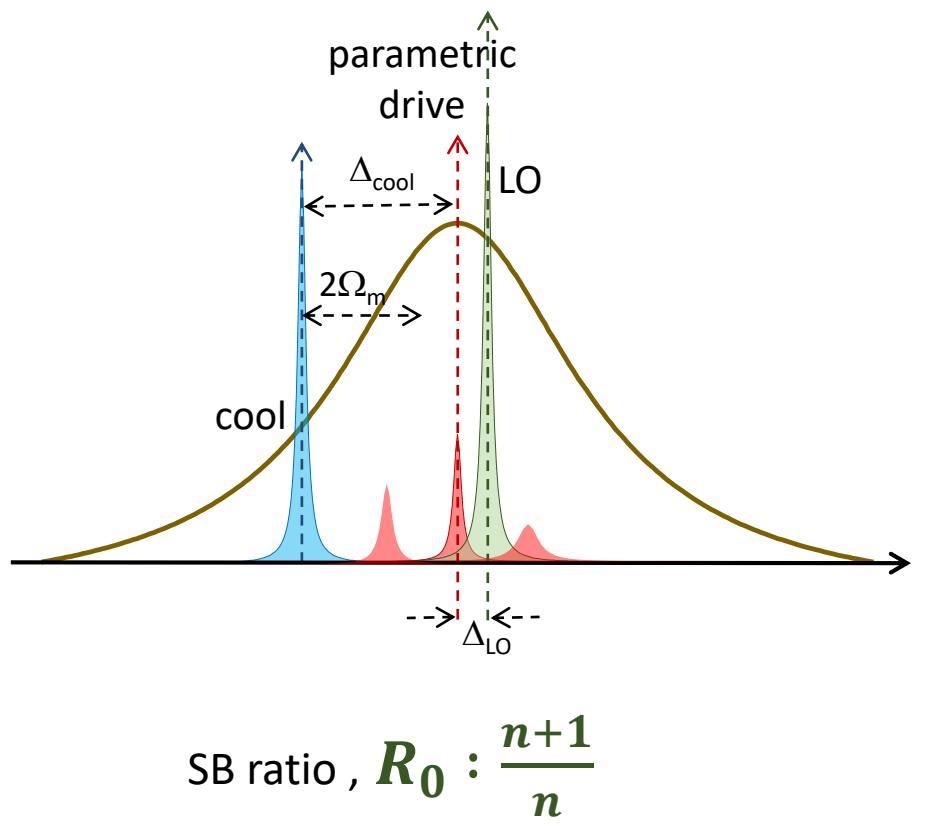
$$I_{pump} = I_{cool} + I_{par}, \text{ where: } I_{par} = \alpha I_{pump}$$

(a) $I_{pump} \uparrow$ keeping ' α ' constant: 's' is constant

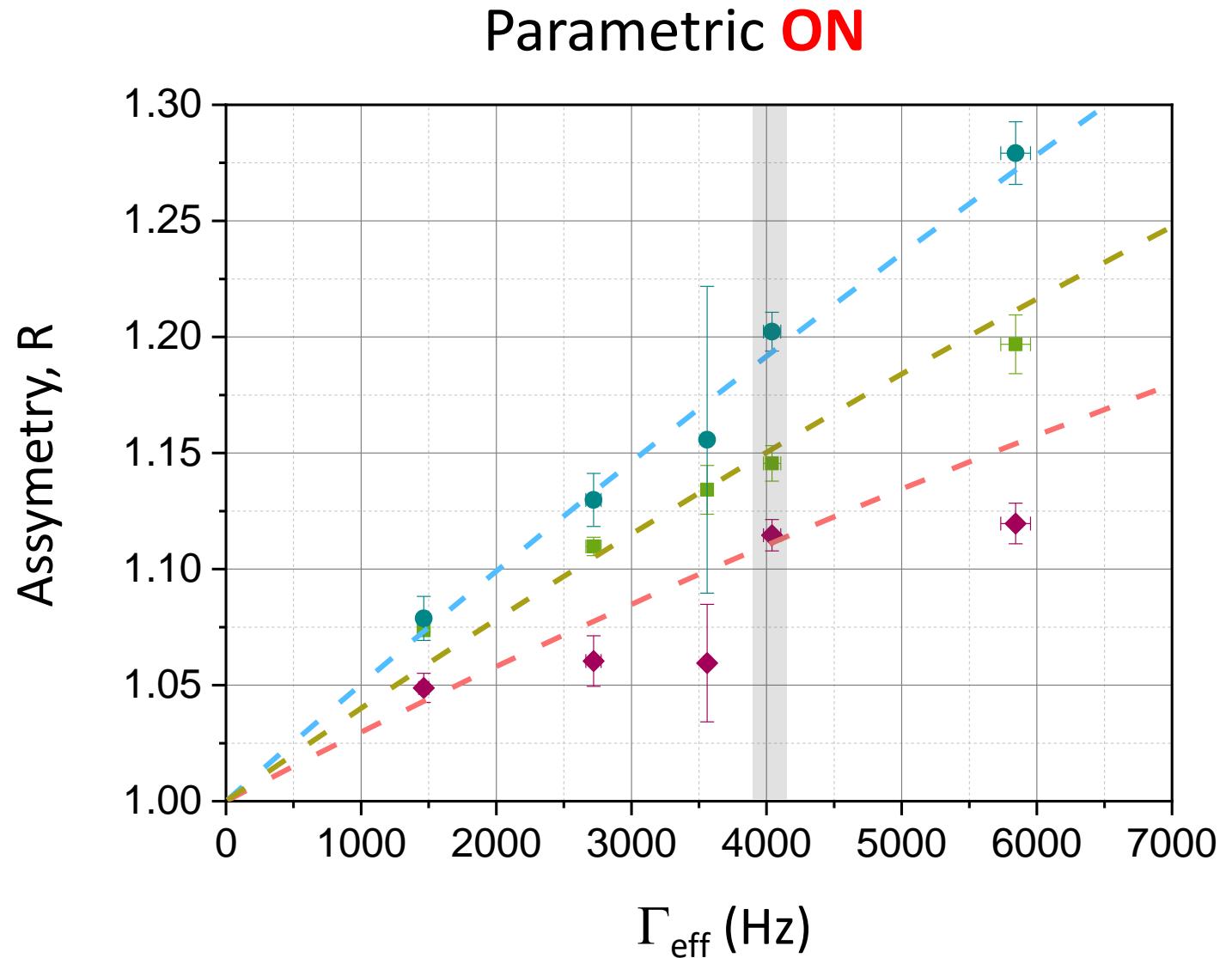
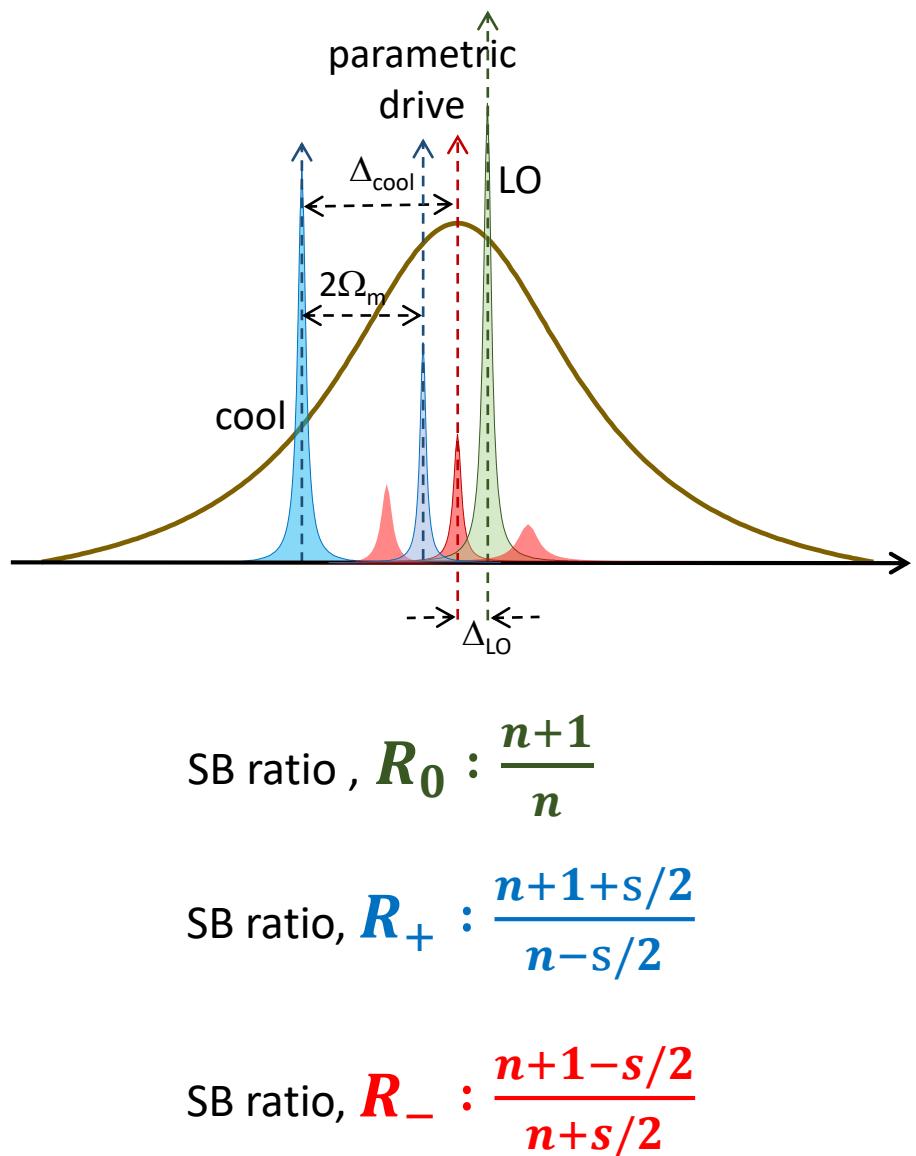
(b) $I_{par} \uparrow$ keeping I_{pump} constant: 's' varies keeping Γ_{eff} constant



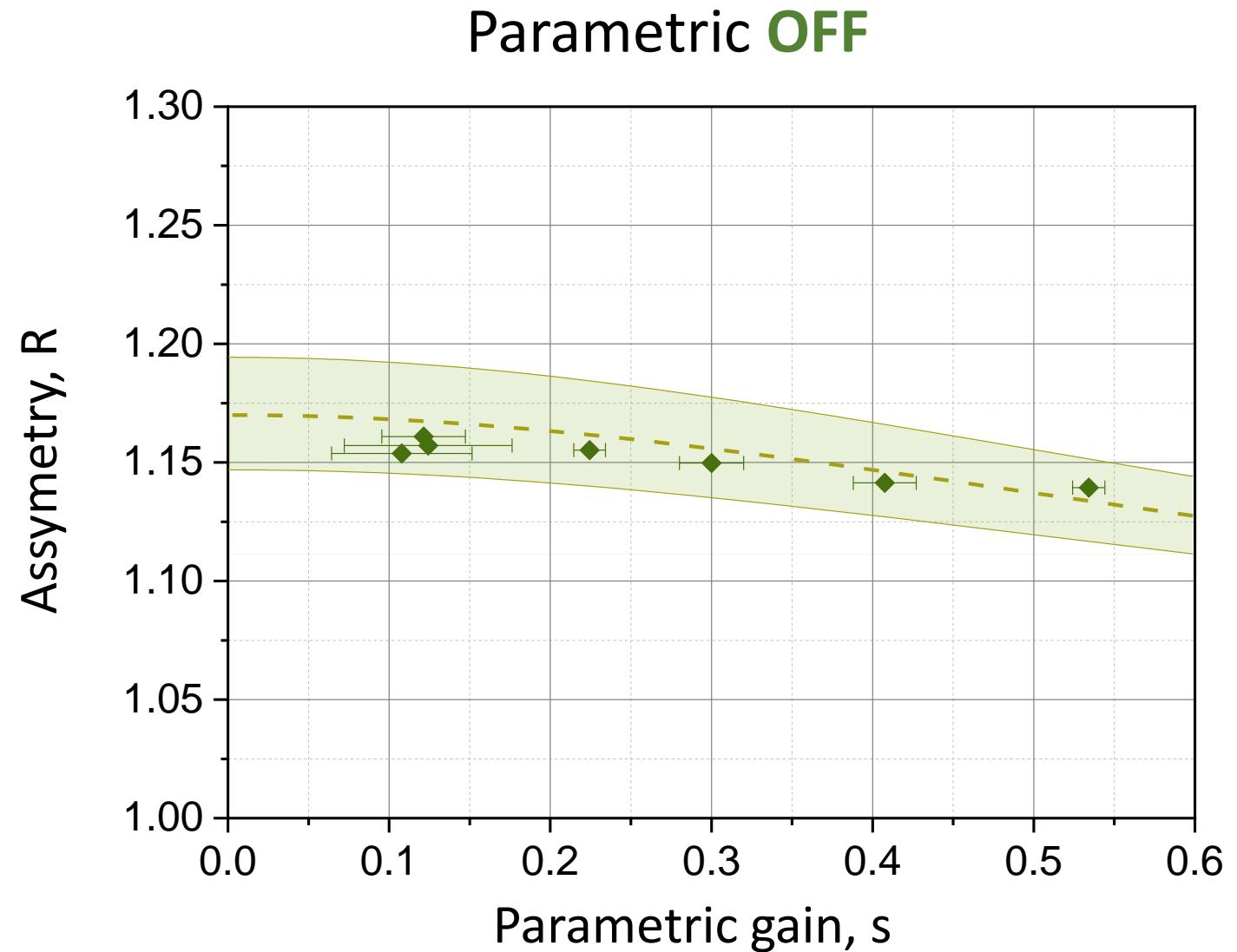
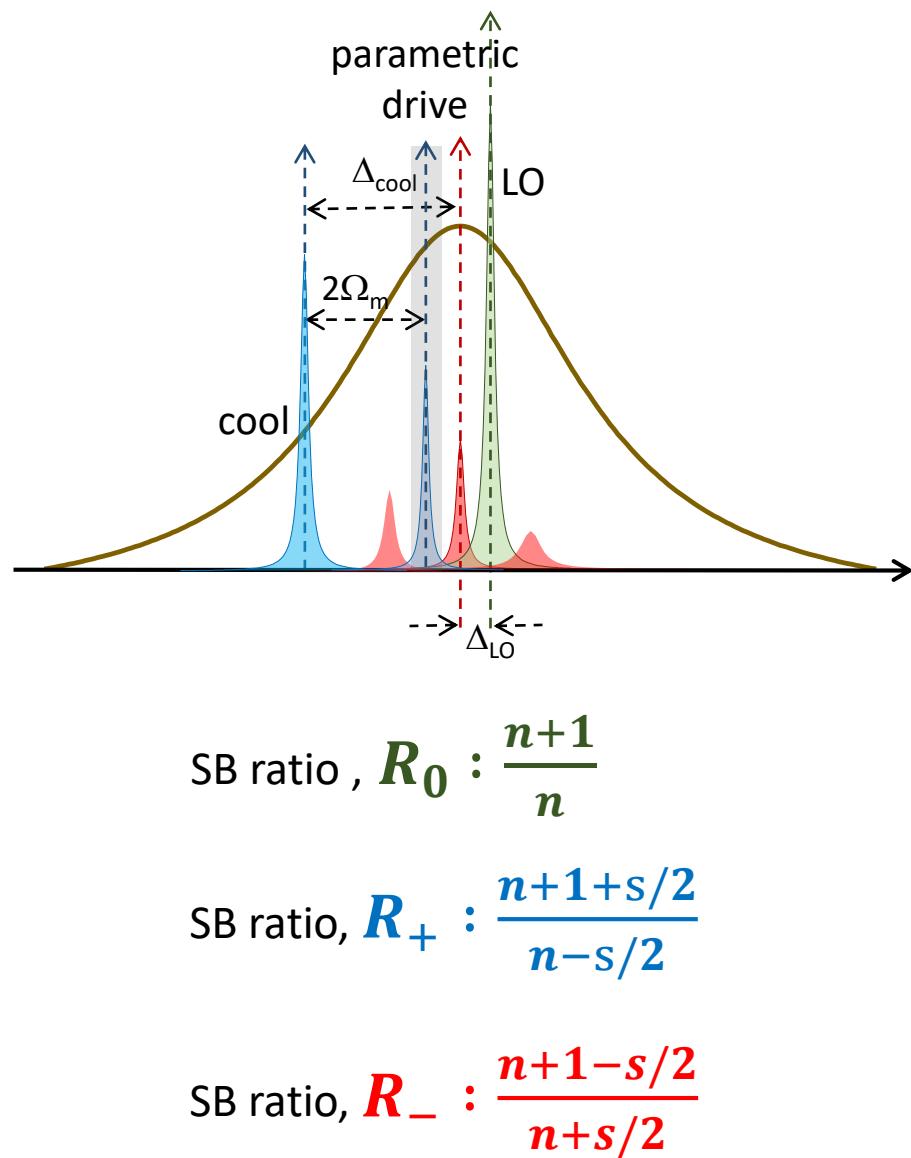
Asymmetry with pump power



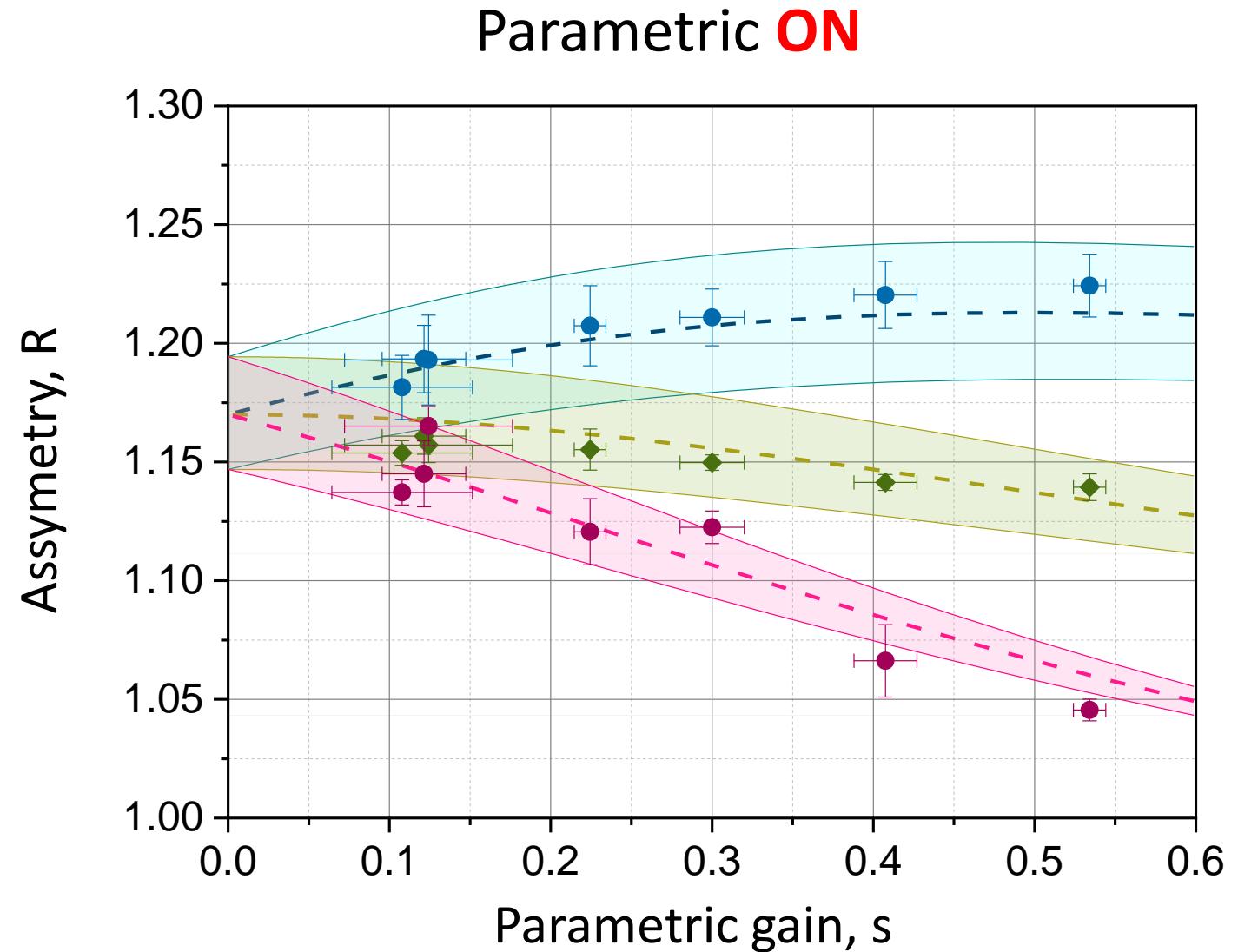
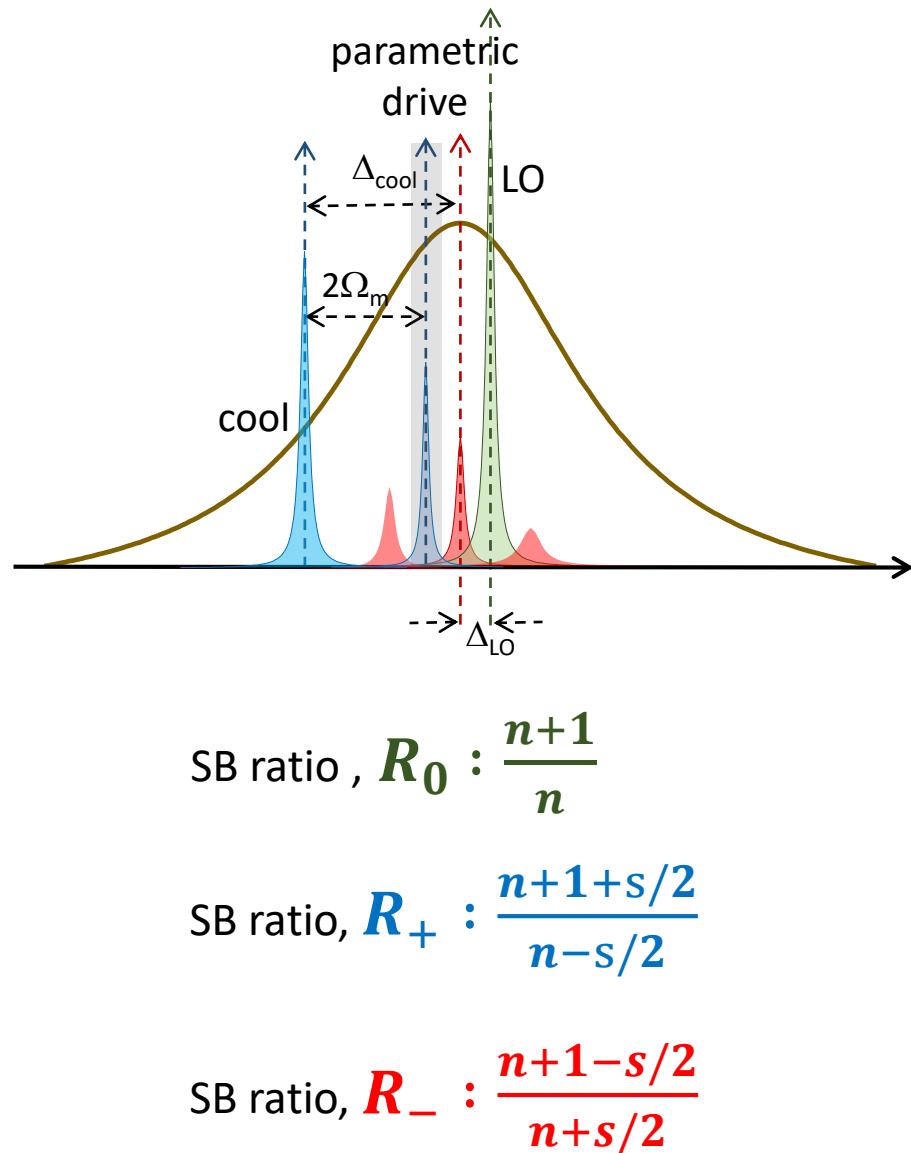
Asymmetry with pump power



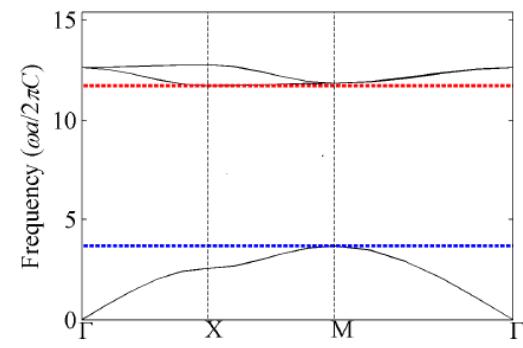
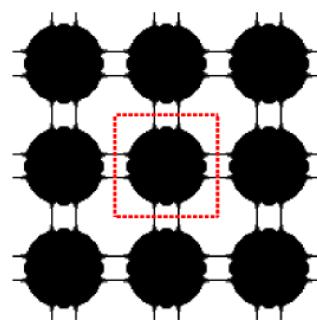
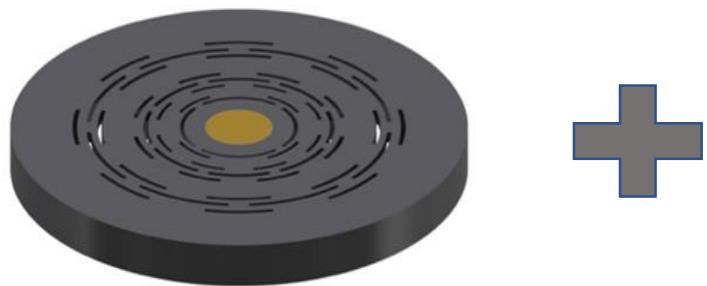
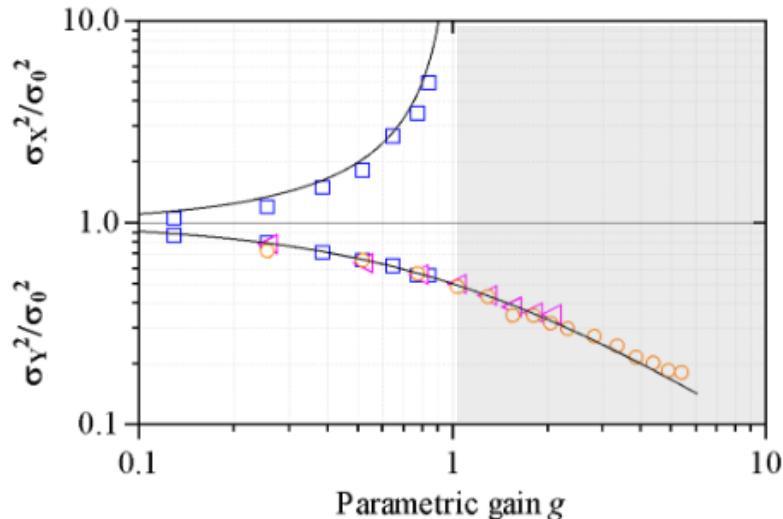
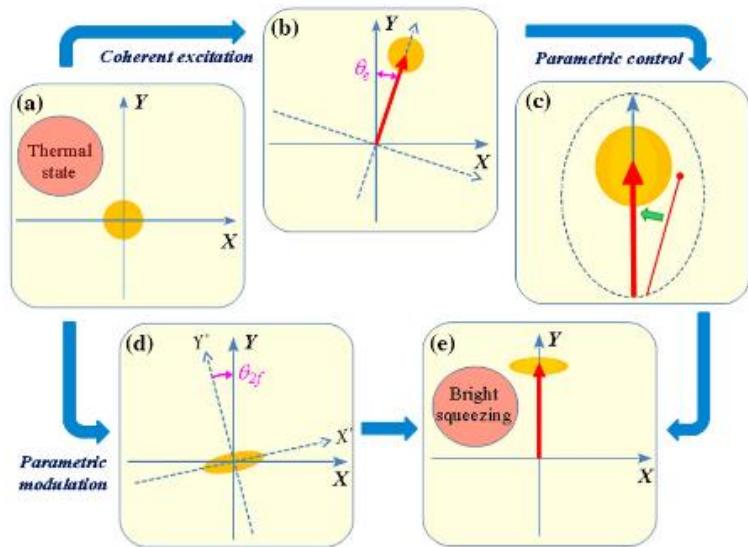
Asymmetry as a function of 's'



Asymmetry as a function of 's'



Perspectives



Parametric feedback control to realize squeezing below 3dB limit close to quantum regime

$Q \sim 10^8$
 $n \sim 10$

Quantum squeezing at room temperature

Thanks



Francesco Marin



Francesco Marino



Paolo Vezio



Massimo Calamai